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A SPECTRUM ANALYSIS
OF SINGLY IONIZED COBALT

By
Nathaniel Ellmaker Hager, Jr.

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A SPECTRUM ANALYSIS OF SINGLY IONIZED COBALT

I. INTRODUCTION

Due in large part to its success in explaining observed spectral regularities, quantum mechanics is today generally accepted as the correct theoretical basis for dealing with atomic and molecular systems. As yet, however, the theory has not yet been developed to the point where the properties of complex systems can be calculated with an accuracy comparable to that of experimental determinations. Experimental spectroscopy continues, as it has in the past, to furnish the best available values for the energy states of particular systems and to provide a proving ground for further development of the theory. The problem of deducing values for the energy levels of particular atoms or molecules from spectral observations is the problem of the spectral analysis.

The analysis of the spectrum of singly ionized cobalt (Co II) was undertaken as the subject of the present thesis at the suggestion of the Spectroscopy Division, National Bureau of Standards. "Of many possibilities, this spectrum was chosen because previous analyses were incomplete, because structures for the spectra of neighboring atoms and ions had been fairly well determined, and because it seemed reasonable to suppose that better measurements could be made.

With regard to previous analyses of Co II, the first regularities were reported by Meggers¹ in 1928. In 1930, Findlay² discovered a few new energy levels, and also made a study of Zeeman patterns which led to a revised classification for several of the levels located by Meggers. Among the new levels found by Findlay were the second members of two series, from which was calculated a value of 17.3 electron volts for the ionization potential. In all, about 200 lines, which determine the relative positions of 55 energy levels, had been identified before the present study was undertaken.

The following account of additions that have been made to the analysis, begins with a brief description of the experimental data employed. Next considered are the important electron configurations and resulting terms to be expected on the basis of theory. The electron configurations are then treated separately and it is pointed out, with supporting evidence, which of the expected terms have actually been found. Appendix I contains a diagram showing the relative positions of the new terms, while values for the individual energy levels are listed in Section VI.

- 1- W. F. Meggers, Journ. Wash. Ac. Sci., 18, 325 (1928)
- 2- J. H. Findlay, Phys. Rev., 36, 5 (1930)

II. OBSERVATIONS

During the course of the present investigation, two new sets of wave length measurements and intensity estimates have been made. Other unpublished experimental data that were available for extending the analysis included wave length measurements in the region from 5000 A to 2000 A made by Dr. Meggers, wave length measurements in the region from 2250 A to 1200 A furnished by Dr. Shenstone, and Zeeman observations made by Dr. Weeks.

One of the new sets of measurements was made from plates furnished by Dr. Kiess of the National Bureau of Standards. These plates cover the wave length region from 5200 A to 2200 A; a condensed spark was used to produce the excitation.

The other set of new measurements was made from spectrograms obtained by the author using a Schuler tube constructed especially for the purpose³. With neon in the tube, the spectrum was photographed from 6400 A to 2000 A with the 21 ft. concave grating spectrograph at Princeton University. Wave length measurements were made in the regions from 6400 A to 5000 A, and from 3000 A to 2000 A. These were the most accurate measurements available, not only because of the large dispersion of the

3- Details of the Schuler tube and accompanying apparatus are given in Appendix II.

spectrograph, but also because of the great sharpness of the lines produced when the Schuler tube is used as a source. It is apparent from examination of the array of known energy levels that these measurements are seldom inaccurate by more than 0.1 reciprocal centimeters. With helium in the Schuler tube, the spectrum was photographed from 2250 Å to 800 Å on the two-meter normal incidence vacuum spectrograph at Princeton, and measurements were made in the wave length region from 1900 Å to 1400 Å.

III. STRUCTURE OF THE SPECTRUM

Part A. of the present section considers the general energy level structure of Co II to be expected on a basis of theory. It contains a list of the terms that could result from each of the important electron configurations and indicates which of these terms have been found. Details of the analysis are discussed in Part B., where the structure of each electron configuration is treated individually.

An energy level diagram is provided in Appendix I. which may be unfolded, and thus made available for continuous reference. This diagram indicates the relative positions of the known terms and it is included primarily to accompany the following discussion.

A. THEORY

The once ionized cobalt atom possesses twenty-six electrons. Eighteen of these are in closed shells resembling the electron structure of argon, leaving eight outer electrons which are active in producing the optical spectrum.

Application of the Aufbauprinzip of Bohr⁴ shows that

4- H. E. White, "Introduction to Atomic Spectra",
p. 79, McGraw-Hill (1934)

the lowest configurations of these eight electrons should be $3d^8$, $3d^7 4s$, and $3d^6 4s^2$. Higher configurations should result from an increase in the quantum numbers n or l associated with one or more of the electrons. Of many possible higher configurations, terms have been found for $3d^7 4p$, $3d^7 5s$ and $3d^6 4s 4p$. The gross structure of the spectrum, as indicated by the relative positions of the various configurations, is shown schematically in the diagram of Appendix I., where the main structural components are enclosed within dashed-line frames.

In Table I. are listed the terms that, according to theory, should arise from each of the electron configurations if it is assumed that there is Russell-Saunders coupling between the electrons. Each Co II term listed is one of a series that converges to a limit which is the Co III term heading the column in the table. For example, the $3d^7 4s \ ^5F$ term of Co II is the lowest member of a series converging to the limit $\ ^4F$, which is a term in the $3d^7$ configuration of Co III.

Most of the terms which have been established are preceded in the table by small letters in accordance with the usual scheme⁵. The unassigned energy levels listed

5- H. N. Russell, A. G. Shenstone, L. A. Turner,
Phys. Rev., 33, 900 (1929)

in Table XIII may belong to those terms which are prefixed with a question mark.

In accordance with Pauli's exclusion principle , the $3d^8$ configuration is expected to give rise to fewer terms than one might otherwise expect. In this case, limit structure cannot be unambiguously assigned to the terms.

Table I.

TERMS ARISING FROM THE ELECTRON
CONFIGURATIONS OF CO II

Co III	3d ⁷	4F	4p	2H	2G	2F	a ² D	2P	b ² D	
Added Electron	4s	a ⁵ F b ³ F	a ⁵ P b ³ P	a ³ H 1 ¹ H	a ³ G 1 ¹ G	3F 1 ¹ F	a ³ D 1 ¹ D	c ³ P 1 ¹ P	3D 1 ¹ D	
	Co II	4p	z ⁵ D ^o z ⁵ F ^o z ⁵ G ^o	z ⁵ S ^o z ⁵ P ^o y ⁵ D ^o	x ³ G ^o z ³ H ^o z ³ I ^o	y ³ F ^o y ³ G ^o z ³ H ^o	? ³ D ^o ? ³ F ^o ? ³ G ^o	? ³ P ^o ? ³ D ^o ? ³ F ^o	? ³ S ^o ? ³ P ^o ? ³ D ^o	? ³ P ^o ? ³ D ^o ? ³ F ^o
z ³ D ^o z ³ F ^o z ³ G ^o			z ³ S ^o ? ³ P ^o y ³ D ^o	1 ¹ G ^o 1 ¹ H ^o 1 ¹ I ^o	1 ¹ F ^o 1 ¹ G ^o 1 ¹ H ^o	1 ¹ D ^o 1 ¹ F ^o 1 ¹ G ^o	1 ¹ P ^o 1 ¹ D ^o 1 ¹ F ^o	1 ¹ S ^o 1 ¹ P ^o 1 ¹ D ^o	1 ¹ P ^o 1 ¹ D ^o 1 ¹ F ^o	
5s			e ⁵ F e ³ F	5P 3P	3H 1H	3G 1G	3F 1F	3D 1D	3P 1P	3D 1D
Co III			3d ⁶ 4s	6D	a ⁴ D	(7 quartets and 15 doublets)				
Co II	4s	a ⁵ D	(quintets, triplets, and singlets based on above limits)							
	4p	7P ^o 7D ^o 7F ^o	5P ^o 5D ^o 5F ^o	(quintets, triplets, and singlets based on above limits)						
		y ⁵ P ^o x ⁵ D ^o y ⁵ F ^o	: ³ P ^o : ³ D ^o : ³ F ^o							
Co III	3d ⁷	(See text on page 7)								
Co II	3d	a ³ F	a ³ P	1S	1D	1G				

? - unassigned levels may be part of term
: - letter not assigned to term

B. DETAILS OF THE ANALYSIS

The low configurations of the spectrum will be discussed individually. Although terms discovered by earlier workers will be mentioned, only the new terms found during this analysis will be treated in detail. Since no singlets have been established, they will be omitted from the following discussion.

THE $3d^8$ CONFIGURATION

As shown in Table I., this configuration should give rise to two triplets. The 3P term was found during the present analysis, while the 3F term was found by Findlay. According to Hund's rule, the 3F term should be lower than the 3P term, and reference to Table XIII. will show that the terms do occur in this order. Since the 3F term provides the lowest series limit for the Co I spectrum⁶, it is not surprising that Findlay found this triplet to be the lowest term in Co II.

6- H. N. Russell, A. S. King, and C. E. Moore,
Phys. Rev., 58, 407 (1940)

Table II⁷ indicates the number of combinations and the number of Zeeman patterns helping to establish each level of the new term a^3P .

Table II.
NEW TERM IN $3d^8$ CONFIGURATION

TERM	J	NUMBER OF COMBINATIONS	NUMBER OF ZEEMAN PATTERNS	ZEEMAN IDENTIFICATION
a^3P	2	21	2	
	1	17	1	3P_1
	0	9	1	3P_0

Evidence from the Zeeman patterns does not conclusively establish the identity of the level which has the highest J-value, because the available patterns are too poorly resolved. But the patterns do not disagree with the identification made by the usual method of wave number differences and intensity rules. (See Table XII)

- 7- If resolved Zeeman patterns provide identification of a level, the last column in the table lists this identification. If no entry appears in the last column, the patterns have not been unambiguously interpreted. The Zeeman patterns are treated more fully in Section V, and the identified combinations are tabulated in Section VII.

THE $3d^7 4s$ CONFIGURATION

Three new triplets have been established and two more have been tentatively established in this group. Table I shows that only one 3F and one 3D have not yet been found, and by means of isoelectronic sequences such as those shown in a later part of this report, it can be shown that these two triplets should be higher in the energy level diagram than the others in this group. This should make them harder to find than the lower terms, because the spectral lines required to establish them may be expected to be much weaker in intensity than those used to establish the lower terms.

Table III⁸ shows the number of combinations and the number of Zeeman patterns helping to establish each level of the new terms. From the table it is apparent that the terms b^3P , c^3F , and a^3D are quite well established because of the large number of combinations, but that the terms a^3G and a^3H are tentative. The term a^3G is especially questionable because some Zeeman patterns do not completely agree with this assignment. This matter is discussed more fully in Section V.

8- For explanation of this table, see footnote 7; this explanation pertains to the succeeding tables of the same type.

Table III.

NEW TERMS IN THE $3d^7 4s$ CONFIGURATION

TERM	J	NUMBER OF COMBINATIONS	NUMBER OF ZEE- MAN PATTERNS	ZEEMAN IDENTI- FICATIONS
b^3P	2	23	3	3P_2
	1	19	2	
	0	12	2	3P_0
c^3P	2	17	3	
	1	18	1	
	0	9	1	
a^3D	3	15	2	3D_3
	2	20	1	3D_2
	1	17		
$a^3G?$	5	4	1	
	4	6	2	3G_4
	3	2		
a^3H	6	4	2	3H_5
	5	6	1	

7 THE 3d 4p CONFIGURATION

Six new terms arising from this configuration have been found during this analysis. Six quintets and three triplets were previously reported by Meggers and Findlay. All of the terms based on the 4F and 4P limits have been established with the possible exception of $3d^7(^4P)4p\ ^3P^0$. Four of the new triplets belong to the group of terms based on the doublet limit structure. (See Table I.)

In addition to the six new terms, sixteen unassigned levels have been found, five of these being located in the region where the above mentioned term, $3d^7(^4P)4p\ ^3P^0$, should be found⁹. It is presumed that the remainder of the sixteen unassigned levels are parts of those terms which have been preceded by a question mark in Table I.

Table IV shows that, of the six new triplets, only $^3D^0$ and $^3S^0$ are well established. The term $y\ ^3G^0$ is especially questionable because two Zeeman patterns do not completely agree with this assignment. This matter is discussed more fully in Section V.

9- The five levels mentioned have been named $^3_2^0$, $^3_3^0$, $^3_4^0$, $^3_5^0$, and $^3_6^0$, and they are listed in Table XIII. As indicated in Table V, they all could be 3P levels, judging from evidence drawn from Zeeman patterns and studies of their selection of levels with which to combine.

Table IV.

NEW TERMS IN THE $3d^7 4p$ CONFIGURATION

TERM	J	NUMBER OF COMBINATIONS	NUMBER OF ZEEMAN PATTERNS	ZEEMAN IDENTIFICATION
y^3F^o	4	5	1	
	3	6	1	
	2	4		
$y^3G^o?$	5	4		
	4	5	1	3G_4
	3	4		
x^3G^o	5	4	1	
	4	3	1	
z^3H^o	6	3	1	3H_6
y^3D^o	3	9	2	
	2	14	2	
	1	10		
z^3S^o	1	11	2	3S_1

Table V lists the sixteen unassigned levels, and it is emphasized that the identifications listed in the last column of the table are quite tentative. It should also be understood that, because of the small number of combinations fixing several of these levels, there is a distinct possibility that some of them may not be real.

Table V.
NEW UNASSIGNED LEVELS

LEVEL	J	NUMBER OF COMBINATIONS	NUMBER OF ZEEMAN PATTERNS	TENTATIVE IDENTIFICATIONS
32^0	0	4		$3p$
33^0	2	8	3	$3p_0$ or $3d_2$
34^0	1	7	2	$3p_2$
35^0	2	9	3	$3p_1$
36^0	1	8	1	$3p_2$
37^0	3	6		$3d_1$
38^0	3	4	1	$3d_3$
39^0	2	7		$3p_3$
310^0	2	3		$3d_2$
311^0	1	2		$3d_2$
312^0	1	5		$3p_1$
313^0	0	3		$3p_1$
314^0	2	5		$3p_0$
315^0	1	5		$3p_2$
316^0	1	7		$3p_1$
317^0	0	2		$3p_1$ 0

There are two reasons for difficulty in grouping these levels into complete terms. First, on the basis of iso-electronic sequence predictions, many terms are expected to fall within a small energy region. This suggests the possibility that the J-levels of two or more terms may be interposed, making it difficult to assign these levels to particular terms on the basis of their relative positions. Second, most of the levels are established by a relatively small number of combinations, and it is consequently difficult to deduce the J and L values of the new levels by ap-

plication of the selection rules.

There are two reasons why the number of combinations establishing the new levels is small. First, most of these levels must belong to terms based on the doublet limit structure, and they combine so weakly with the established terms based on the quartet limit structure that many of the satellites are missing. Second, many of the lines essential to performing this phase of the analysis fall in a wavelength region where they are likely to be hidden by lines of the arc spectrum. Both of these difficulties are illustrated in Table VI, where it is shown how the even term a^3D , which is based on the doublet limit structure, combines with the odd terms z^3F^0 , z^3D^0 , and y^3D^0 , which are based on the quartet limit structure. In this case, because of other combinations, there can be little doubt about the assignments.

Table VI.

WEAK INTER-SYSTEM COMBINATIONS

		z^3F^0			z^3D^0			y^3D^0		
		4	3	2	3	2	1	3	2	1
	1			1		-	10		0	0
a^3D	2		2	2	3	15	CoI	0	3	2
	3	10	CoI	-	50	-		10	-	

The intensities in Table VI were estimated according to an arbitrary scale on which 200 represents the intensity of the brightest lines in the spectrum and 0 represents the intensity of lines which are so weak that they can just barely be observed. It is apparent that the above combinations are weak, and it is not surprising that some of the satellites are missing. Two of the satellites coincide with and may be obscured by arc lines. In the wave length region from 5000 Å to 3000 Å, where the lines arising from these combinations are found, the plates are so blackened with lines that it is certain that many of them are hidden by or blended with others.

For the same reasons it is difficult to establish combinations between the odd terms of the doublet limit system and the even terms b^3F and b^3P of the quartet limit system. These odd terms also combine poorly with the even terms a^3P and a^3F and the $3d^8$ configuration.

THE $3d^7 5s$ CONFIGURATION

The second series members e^3F and e^5F were established by Findlay, the first series members being the terms $3d^7(^4F)4s\ b^3F$ and $3d^7(^4F)4s\ a^5F$ respectively.

It is in connection with higher series members such as these that the Schuler tube source should prove particularly useful. The new plates did, in fact, reveal 14 lines not observed by Findlay, representing transitions from e^3F and e^5F to lower terms. Not only was the identity of these terms more firmly established by the combinations involving the new lines, but also the energy values of the levels were determined more accurately because the Schuler tube yields sharper lines than the spark source used by Findlay. The lines produced by the condensed spark are known to be not only very broad and fuzzy, but also to be asymmetric in such a way as to introduce at times a consistent error in wave length measurements. Table VII compares Findlay's term values with the revised term values.

Table VII.
ADJUSTMENT OF FINDLAY'S ENERGY VALUES

TERM	J	FINDLAY'S VALUES C_m^{-1}	REVISED VALUES C_m^{-1}	DIFFERENCE C_m^{-1}
e^5F	5	84012.3	84013.70	1.4
	4	84584.8	84586.35	1.6
	3	85165.3	85167.10	1.8
	2	85593.9	85595.70	1.8
	1	85874.1	85875.84	1.7
e^3F	4	85479.2	85481.07	1.9
	3	86343.8	86346.19	2.4
	2	86937.7	86939.60	1.9

Attempts at finding third members of these series, using the wave length measurements from the Schuler tube, were not successful. Even the second series members of the two series whose lowest members are b^3P and a^5P were not established. This was disappointing in view of the expected enhancement in the Schuler tube of lines associated with these upper terms. A search was also made to locate terms in the $3d^7 4d$ configuration, but here again there was no success.

THE $3d^6 4s^2$ CONFIGURATION

Although this is the lowest configuration in Fe I, it is quite high in Co II, and prior to this analysis none of its levels were known. Of the terms arising from the configuration, only the 5D has been found. According to Table I, there should also be triplets, but they are predicted to be higher than $80,000\text{ cm}^{-1}$ above the normal state, and have not been established.

Predictions (made by the method described in Section IV) indicate that the 5D term should fall within a few thousand wave numbers of $37,500\text{ cm}^{-1}$. The new term was actually found to be at $41,000\text{ cm}^{-1}$. The new wave length measurements needed for establishing this term were taken from Kiess' plates in the wave length region above 5000 Å.

Table VIII shows the number of combinations establishing each level of the new term. No Zeeman data were available in the wave length region above 5000 A.

Table VIII
NEW TERM IN THE $3d^6 4s^2$ CONFIGURATION

TERM	J	NUMBER OF COMBINATIONS
a^5D	4	10
	3	14
	2	11
	1	7

THE $3d^6 4s4p$ CONFIGURATION

There should be a large number of septets, quintets, triplets and singlets in this configuration, but space has been taken in Table I for listing only a small group of these terms because so few of them have been established. The terms which have been established are listed in Table IX.

Table IX.

THE NEW TERMS IN THE $3d^6 4s4p$ CONFIGURATION

TERM	J	NUMBER OF COMBINATIONS
x^5D^o	4	5
	3	4
	2	4
	1	3
	0	1
y^5F^o	5	3
	4	3
	3	3
	2	2
y^5P^o	3	3
	2	3
	1	2
3^oD	3	2
	2	3
3^oF	4	2
	3	2
	2	2
3^oP	2	3
	1	3

Although the table shows that the levels are not well established as far as the number of combinations is concerned, confidence in the reality of both the quintet and the triplet terms may be gained by examination of the line list in the vacuum region, where the lines used in establishing these terms are found. In this region

there is a large average wave length interval between lines, and an unusually large percentage of these lines arise from combinations involving the new multiplets.¹⁰ The importance of this is that the paucity of lines reduces the chance of finding by coincidence the proper wave number differences for establishing terms.

UNASSIGNED TERMS

Table X indicates the number of combinations helping to establish some terms which were found near $100,000 \text{ cm}^{-1}$ above the ground state. These terms could belong to the $3d^6 4s 4p$, $3d^7 4f$, or $3d^7 5p$ configurations, but it is not possible at present to make a configuration assignment.

TABLE X.
NEW UNASSIGNED TERMS

TERM	J	NUMBER OF COMBINATIONS
$x^5 F^0$	5	2
	4	4
	3	4
	2	4
$v^5 D^0$	4	6
	3	5
	2	2
$w^5 D^0$	4	3
	3	3
	2	2

10- See Appendix III.

IV. COMPARISONS WITH NEIGHBORING SPECTRA

A. ISOELECTRONIC SPECTRA

If, by ionization, one electron is removed from the cobalt atom and two are removed from the nickel atom, then both have the same number of electrons as neutral iron, and the three elements form an isoelectronic sequence. Because they have the same electron configurations, the three atoms should give rise to similar spectra and systems of energy levels.

This similarity affords a method of using data taken from the analyses of Fe I¹¹ and Ni III¹² both for predicting the energy values of unknown Co II terms relative to the established system of levels, and for substantiating the identity of terms which have already been established.

Use of the familiar Mosley diagram for studying isoelectronic sequences is suggested by Hertz's relation

$$T = \frac{R(Z - \sigma)^2}{n^2} \quad (1)$$

where T is called the term value, and is the energy difference between a term in a given series and the limit of that series. R is the Rydberg constant, Z is the atomic number, n is the principle quantum number of the valence

11- H. N. Russell and C. E. Moore, Trans. Am. Phil. Soc., 34, 113 (1944)

12- A. G. Shenstone, unpublished material (1950)

electron, and σ is the so-called screening constant, which accounts for the reduction of the effectiveness of the nuclear charge due to the presence of all of the electrons other than the valence electron. It is possibly misleading that σ is called a constant because it has different values for the different terms in the spectrum of one atom, and it has different values for corresponding terms in different atoms.

Transposing equation (i) and taking the square root of both sides gives the expression

$$\sqrt{\frac{T}{R}} = \frac{Z - \sigma}{n} \quad (11)$$

To make a Mosley diagram, values of $\sqrt{T/R}$ are plotted as ordinates and values of Z are plotted as abscissae. According to equation (11) such a diagram may produce straight lines, provided that the behavior of σ does not upset the linear relationship between the variables.

Gibbs and White¹³ have made an extensive study of isoelectronic sequences in complex spectra, and their Mosley diagrams show lines which are not only quite straight, but also, for those terms having the same value of n , the lines are parallel with each other. The parallelism requires that the difference between values of σ for two

13- R. C. Gibbs and H. E. White, Phys. Rev., 29, 426 (1927), Phys. Rev., 29, 655 (1927), and Proc. Nat. Acad. Sci., 13, 525 (1927)

terms be the same for all elements in the sequence, and the straightness of the lines requires that if σ varies with Z for a given term, this variation must be linear with respect to changes in Z .

The energy difference between two terms belonging to the same element and having the same value of n may be written

$$\begin{aligned} T_1 - T_2 &= \frac{R}{n^2} [(Z - \sigma_1)^2 - (Z - \sigma_2)^2] \\ &= \frac{R(\sigma_1 - \sigma_2)}{n^2} [Z + (\sigma_1 + \sigma_2)] \end{aligned}$$

If $(\sigma_1 - \sigma_2)$ is constant and if $(\sigma_1 + \sigma_2)$ is a linear function of Z , as suggested by the parallel and straight lines on the Mosley diagram, then we may write

$$T_1 - T_2 = AZ + B \quad (111)$$

where A and B are constants for a given value of n , which do not depend on Z .

Therefore, as a substitute for plotting the Mosley diagram, graphs of the type shown in Figs. 1 and 2 may be used. The particular advantage of this type of plot is that it may be made without knowledge of the position of the series limits, and it is used here because the limit structures of Co II and Ni III are not yet well established.

In the diagrams shown, the term which lies lowest in the Co II spectrum has been arbitrarily assigned an energy level value of zero in all three spectra, and the energy differences between this term and the other terms in the same configuration have been plotted with respect to the same scale for all three elements.

Since the individual J-levels of each term would not be resolved in drawings made to the scale shown in the figures, a single point has been plotted to represent the "center of gravity" of the terms. The energy value of the center of gravity of a term having more than one J-level is computed from the equation

$$E_c = \frac{\sum_J E_J (2J+1)}{\sum_J (2J+1)}$$

where E_c is the energy value of the center of gravity and E_J is the energy value of an individual J-level.

Although all of the new terms found in Co II except $3d^8 \ ^3P$ have been found in Fe I, only a few have been reported in the case of Ni III. Fig. 1 shows that straight lines are obtained for the two new terms $3d^7(^2G)4s \ ^3G$, and $3d^7(^4P)4s \ ^3P$. Fig. 2 shows a straight line for the term $3d^7(^2G)4p \ ^3G^o$, and a line for the term $3d^7(^2G)4p \ ^3F^o$ which is no less straight than are the lines for several of the better established terms.

Fig. 1

ISOELECTRONIC SEQUENCE DIAGRAM FOR EVEN TERMS

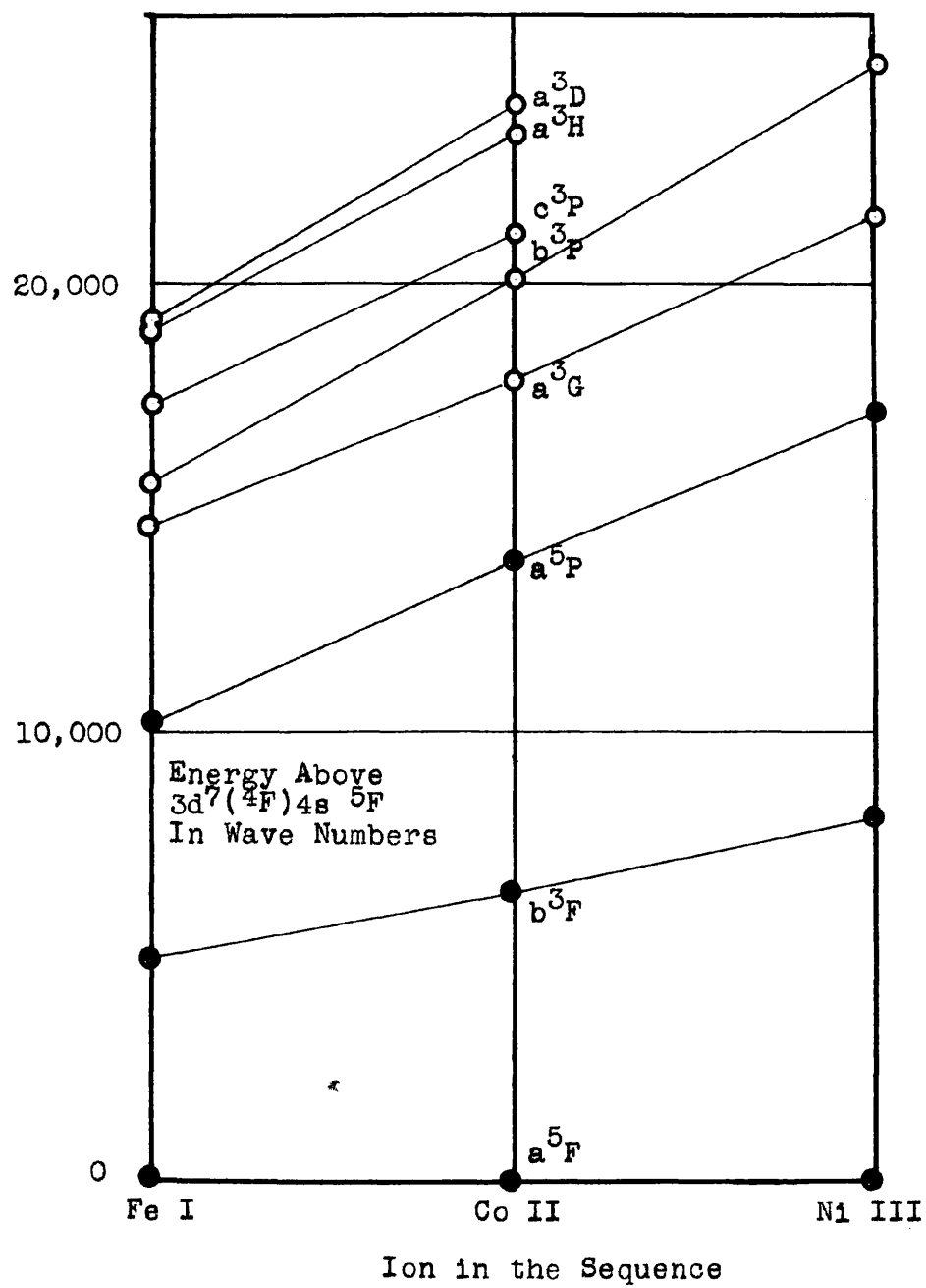
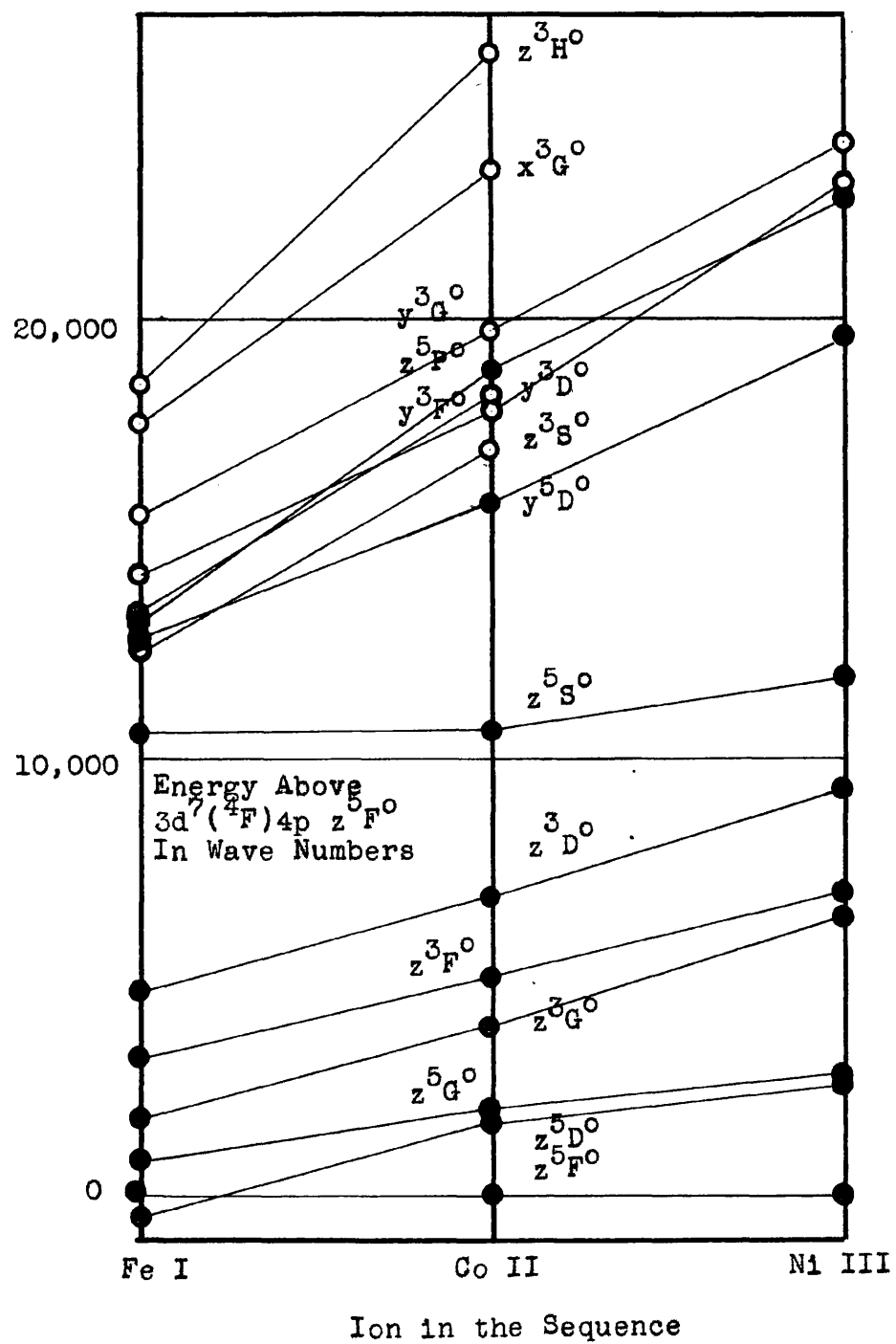


Fig. 2

ISOELECTRONIC SEQUENCE DIAGRAM FOR ODD TERMS



The isoelectronic sequence drawings do not positively confirm or deny the correctness of the term identifications, but it is encouraging to note that terms in the three spectra are located as expected from Hertz's formula and from implications drawn from the isoelectronic studies of Gibbs and White.

Additional similarities may be shown to exist between the term of Fe I, Co II, and Ni III by comparing the intervals between the levels of the same term in each of the three spectra. Table XI is devoted to such a comparison. The intervals between the levels of each term have been tabulated for the three spectra, and two other columns indicate the ratios of the intervals in one spectrum to those in the spectrum of the neighboring atom in the isoelectronic sequence. The comparisons show striking consistencies in many cases, and this material is included because it helps confirm the identity of terms which have been established in the $3d^6 4s 4p$, $3d^6 4s^2$, $3d^7 5s$, and $3d^7 4s$ groups.

There are discrepancies in the case of the terms in the $3d^7 4p$ configuration, but these irregularities may be caused by perturbations arising from the fact that there are so many terms crowded into such a small energy region.

Although the nature of this perturbation is discussed by White¹⁴, it will be briefly stated here that if two levels, belonging to the same or to different configurations, have approximately the same energy value, the same parity, and the same J value, then there may be an apparent repulsion between the two levels. The displacement of the two levels will cause a rearrangement of the intervals in both terms, and hence lead to irregularities of the type seen in Table XI.

As an example, the intervals in the terms $3d^7(^4F)4p\ ^5F^0$, $3d^7(^4F)4p\ ^5G^0$, and $3d^7(^4P)4p\ ^5D^0$ are subject to this type of perturbation. Since the terms of the $3d^74p$ and the $3d^64s4p$ configurations of Fe I fall in the same energy level region and since both configurations include a very large number of terms, it seems reasonable that the repulsion effect should cause irregularities when the intervals within the terms in this spectrum are compared with those within the terms of the Co II spectrum, where the $3d^74p$ and the $3d^64s4p$ configurations do not fall in the same energy region.

14- H. E. White, op. cit., p. 390

Table XI.
MULTIPLYET INTERVAL COMPARISONS

CONFIG.	DESIG.	J	FE I DIFF.	RATIO	CO II DIFF.	RATIO	NI III DIFF.
$3d^8$	a^3F	4	539.1	1.8	950.28	1.4	1360.7
		3	352.6	1.8	647.15	1.4	908.9
		2					
$3d^7(^4F)4s$	a^5F	5	448.5	1.5	678.53	1.4	954.2
		4	351.3	1.5	531.84	1.4	748.2
		3	257.6	1.5	389.25	1.4	545.9
		2	168.9	1.5	254.61	1.4	356.2
		1					
$3d^7(^4F)4s$	b^3F	4	584.7	1.5	895.58	1.4	1267.0
		3	407.6	1.5	613.53	1.4	866.3
		2					
$3d^8$	a^3P	2			143.66	2.2	316.2
		1			188.93	1.3	252.9
		0					
$3d^7(^4P)4s$	a^5P	3	176.8	1.5	259.9	1.2	260.0
		2	200.4	1.5	307.2	1.5	307.0
		1					
$3d^7(^2G)4s$	a^3G	5	283.4	1.4	384.7?	1.4	523.1
		4	250.3	0.8	217.3?	2.7	590.9
		3					

Table XI. (Continued)

CONFIG.	DESIG.	J	FE I DIFF.	RATIO	CO II DIFF.	RATIO	NI III DIFF.
$3d^7(4p)4s$	b^3P	2					
		1	108.5	1.8	193.1		
		0	104.9	1.4	143.8		
$3d^7(2p)4s$	c^3P	2					
		1	436.3	1.0	431.4		
		0	319.6	1.7	543.8		
$3d^7(2H)4s$	a^3H	6					
		5	245.1	1.5	363.4		
		4	276.6				
$3d^7(2D)4s$	a^3D	3					
		2	398.7	1.6	627.4		
		1	-217.2	2.4	-526.5		
$3d^6 4s^2$	a^5D	4					
		3	415.9	1.5	619.0		
		2	288.1	1.5	424.0		
		1	184.1	1.5	270.8		
		0	89.9				

Table XI. (Continued)

CONFIG.	DESIG.	J	FE I DIFF.	RATIO	CO II DIFF.	RATIO	NI III DIFF.
$3d^7(^4F)5s$	e^5F	5					
		4	372.5	1.5	572.65		
		3	377.6	1.5	580.75		
		2	281.1	1.5	428.60		
		1	184.7	1.5	280.14		
$3d^7(^4F)5s$	e^3F	4	570.9	1.5	865.12		
		3	396.5	1.5	593.41		
		2					
$3d^7(^4F)4p$	z^5F^o	5					
		4	344.1	0.5	181.06	0.9	158.7
		3	289.2	2.0	593.23	1.4	849.8
		2	218.5	2.2	480.76	1.4	693.2
		1	144.9	2.3	333.64	1.4	487.4
$3d^7(^4F)4p$	z^5D^o	4					
		3	411.2	1.7	718.30	1.4	1036.8
		2	294.4	1.7	498.36	1.4	715.7
		1	215.5	1.4	311.39	1.4	444.4
		0	104.5	1.4	146.77	1.4	199.8

Table XI. (Continued)

CONFIG.	DESIG.	J	FE I DIFF.	RATIO	CO II DIFF.	RATIO	NI III DIFF.
$3d^7(^4F)4p$	z^5G^o	6					
			-61.5	-4.4	267.29	1.3	353.4
		5					
			474.9	1.0	461.66	1.2	564.2
		4					
			354.3	1.0	343.40	1.2	404.9
$3d^7(^4F)4p$	z^3G^o	3					
			244.8	1.0	237.43	1.1	260.9
		2					
		5					
			388.4	2.0	792.35	1.8	1402.0
$3d^7(^4F)4p$	z^3F^o	4					
			476.6	1.6	684.01	1.6	1058.9
		3					
			358.4	1.5	532.73	1.6	864.3
		2					
$3d^7(^4F)4p$	z^3D^o	3					
			502.7	1.4	717.49	1.3	924.3
		2					
			317.7	1.4	454.88	1.3	602.7
		1					
$3d^7(^4P)4p$	y^5D^o	4					
			24.2	-6.1	-147.6	2.4	-358.4
		3					
			143.5	0.1	19.7	-2.0	-41.1
		2					
			288.7	0.3	89.9	0.5	44.8
$3d^7(^4P)4p$	y^5D^o	1					
			-5.73	-18.7	107.6		
$3d^7(^4P)4p$	y^5D^o	0					

Table XI. (Continued)

CONFIG.	DESIG.	J	FE I DIFF.	RATIO	CO II DIFF.	RATIO	NI III DIFF.
$3d^7(^4P)4p$	z^5P^o	3	196.9	0.1	23.0	4.8	109.4
		2	126.4	2.4	298.0		
		1					
$3d^7(^4P)4p$	y^3D^o	3	118.9	0.2	29.1		
		2	136.0	1.8	249.2		
		1					
$3d^7(^2G)4p$	y^3F^o	4	203.6	4.2	850.0	1.6	1365.9
		3	104.3	6.0	657.0	1.6	1074.4
		2					
$3d^7(^2G)4p$	y^3G^o	5	167.0	4.5	752.4?	1.0	723.6
		4	222.7	0.1	20.6?		
		3					
$3d^7(^2H)4p$	x^3G^o	5	294.3	2.2	640.4?		
		4	157.6				
		3					
$3d^64s(^6D)4p$	x^5D^o	4	240.2	1.0	243.7		
		3	199.5	1.2	245.7		
		2	139.7	1.3	182.7		
		1	71.1	1.3	95.4?		
		0					

Table XI. (Continued)

CONFIG.	DESIG.	J	FE I DIFF.	RATIO	CO II DIFF.	RATIO	NI III DIFF.
$3d^6 4s(^6D)4p$	y^5F^o	5					
			292.3	1.3	373.1		
		4					
			227.9	1.3	284.3		
		3					
			164.9	2.0	332.0		
		2					
			106.7				
		1					
$3d^6 4s(^6D)4p$	y^5P^o	3					
			412.7	1.6	668.4		
		2					
			263.7	1.6	421.7		
		1					
$3d^6 4s(a^4D)4p$	$^3D^o$	3					
			363.7	1.4	508.7		
		2					
			251.0				
		1					
$3d^6 4s(a^4D)4p$	$^3F^o$	4					
			497.8	1.6	784.7		
		3					
			328.9	1.6	521.3		
		2					
$3d^6 4s(a^4D)4p$	$^3P^o$	2					
			415.9	1.4	586.8		
		1					
			192.8				
		0					

B. HOMOLOGOUS SPECTRA

Those elements of the first long period which are immediate neighbors of cobalt in the periodic table give rise, when singly ionized, to the configurations $3d^n$, $3d^{n-1}4s$, and $3d^{n-2}4s^2$. These elements are said to form a homologous sequence.¹⁵

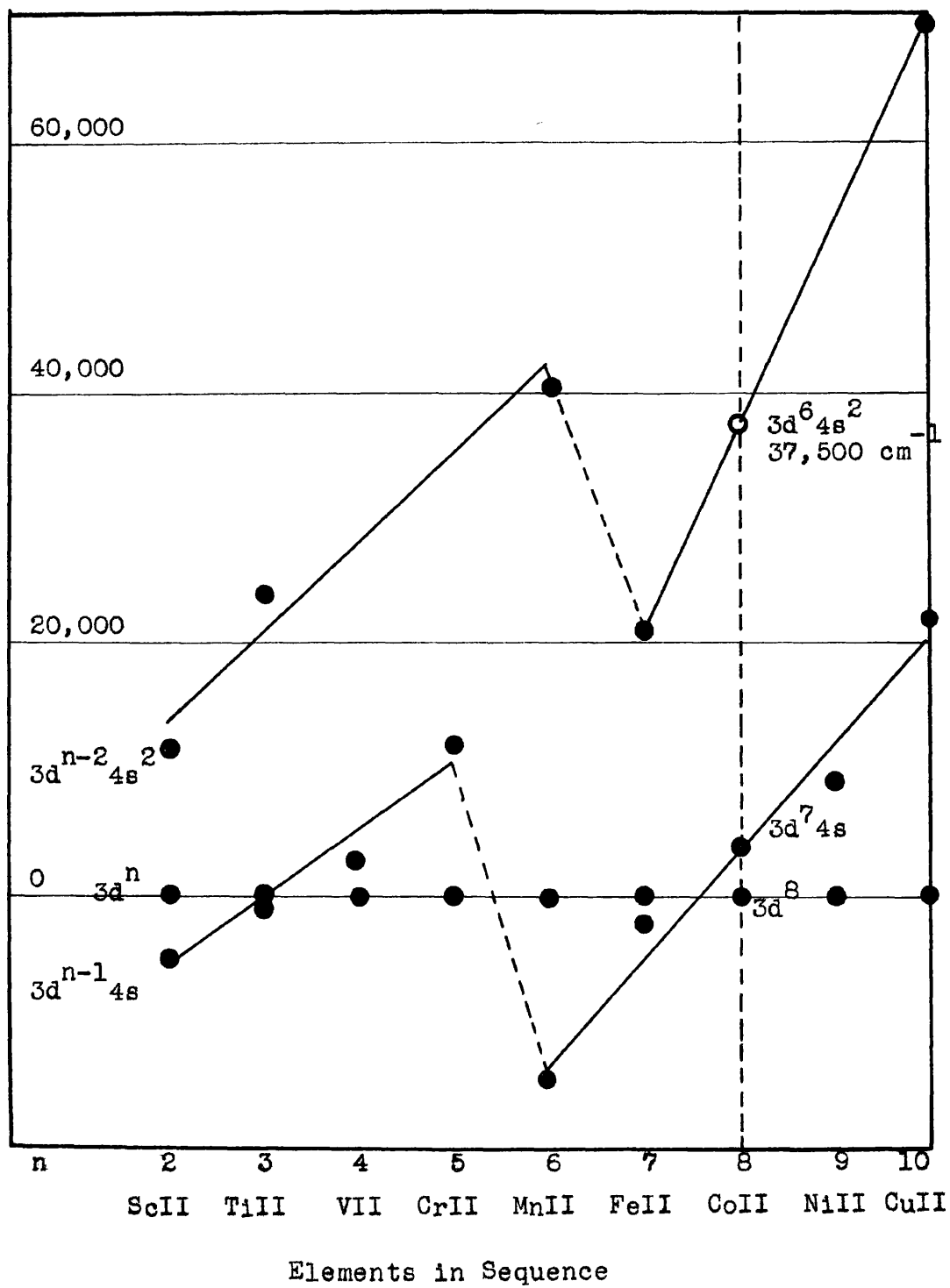
If, for each member of this sequence, the lowest lying level in each configuration is plotted with respect to the lowest level in the $3d^n$ configuration of the same element, as shown in Fig. 3¹⁶, there is sufficient regularity to provide a basis for predicting the approximate position of the $3d^6 4s^2$ configuration of Co II.

In the figure it is seen that this configuration should fall at an energy level which is higher than the level of the $3d^8$ configuration by about 37,000 wave numbers, with a possible error of a few thousand wave numbers.

15- For further illustration see H. E. White, op. cit., pp. 277-78.

16- Data for other elements obtained from R. F. Bacher and S. Goudsmit, "Atomic Energy States", McGraw-Hill (1932).

Fig. 3
HOMOLOGOUS SEQUENCE



V. ZEEMAN PATTERNS

Weeks has observed Zeeman patterns for some of the newly identified lines. In Table XII, data from these observed patterns are compared with values calculated from the Lande theory. In cases where no assignment has been made to one of the two energy levels involved in a combination, no comparison with theory is possible, and only the observed data have been tabulated.

In the third column of the table, "O" designates the experimentally observed pattern, and "C" designates the pattern calculated from theory. The numbers enclosed within parentheses represent measurements of the intervals between the components of the pattern obtained when the observations are made along a line parallel with the magnetic field, and the numbers outside the parentheses represent the measurements obtained when the observations are made along a line perpendicular to the field.

There are several cases in which the observed pattern does not fully agree with that expected from theory. Because the plates were not available for further examination, it was not possible to reconcile these discrepancies completely, and the remarks in the last column of the table are necessarily speculative.

Table XII.
ZEEMAN PATTERN COMPARISONS

WAVELENGTH	COMBINATION	ZEEMAN PATTERN	REMARKS
4160.69 A.	$a^3D_3 - z^3D_3^0$	O(0.00)1.32 C(0.00)1.33	
4145.15	$a^3D_2 - z^3D_2^0$	O(0.00)1.17 C(0.00)1.17	
3754.71	$b^3P_2 - z^3D_3^0$	O(0.02, 0.16, -)0.90, 1.07, 1.25, - , - C(0.00, 0.17, 0.3)1.00, 1.17, 1.33, 1.50, 1.67	Missing components should be weak according to selection rules
3553.93	$b^3P_0 - z^3D_1^0$	O(0.00)0.50 C(0.00)0.50	
2636.04	$a^3D_3 - ^3S_2^0$	O(0.00)1.11	
2628.80	$b^3P_0 - z^3S_1^0$	O(0.00)2.04 C(0.00)2.00	
2613.51	$a^3P_2 - z^3D_3^0$	O(0.00, - , -)1.28 C(0.00, 0.17, 0.33)1.00, 1.17 1.33, 1.50, 1.67	Weak σ components missing; π com- ponents may be unresolved
2605.71	$b^3P_1 - z^3S_1^0$	O(0.53)0.95, 1.49, 1.96 C(0.50)1.00, 1.50, 2.00	
2583.18	$b^3P_2 - y^3D_3^0$	O(0.00, - , -)1.16 C(0.00, 0.17, 0.33)1.00, 1.17 1.33, 1.50, 1.67	Weak σ components missing; π com- ponents may be unresolved

Table XII. (Continued)

WAVELENGTH	COMBINATION	ZEEMAN PATTERN	REMARKS
2574.90 A.	$a^3P_1 - z^3D_2^0$	O(0.01, 0.35)0.82, 1.16, 1.49 C(0.00, 0.33)0.83, 1.17, 1.50	
2565.39	$a^3P_2 - z^3D_2^0$	O(0.45)0.95, 1.17, 1.38, 1.59 C(0.33, 0.67)0.83, 1.17, 1.50, 1.83	pattern may be partly un- resolved; discrepancies unexplained
2557.36	$a^3P_0 - z^3D_1^0$	O(0.00)0.49 C(0.00)0.50	
2552.40	$c^3P_0 - z^3D_1^0$	O(0.00)1.49	
2546.74	$b^3P_0 - z^5P_1^0$	O(0.27)1.04 C(0.00)2.50	Zeeman effect indicates identification incorrect
2540.63	$b^3P_1 - y^3D_2^0$	O(0.00, -)0.82, -, - C(0.00, 0.33)0.83, 1.17, 1.50	Weak compon- ents missing
2530.09	$b^3P_2 - y^3D_3^0$	O(- , - , -) 1.14 C(0.00, 0.17, 0.33)1.00, 1.17 1.33, 1.50, 1.67	π components too weak; σ components may be un- resolved
2528.21	$b^3P_2 - y^3D_2^0$	(- , 0.37) 1.30?, 1.47? C(0.33, 0.67)0.83, 1.17, 1.50, 1.83	Discrepancies unexplained

Table XII. (Continued)

WAVE LENGTH	COMBINATION	ZEEMAN PATTERN	REMARKS
2497.48	$c^3P_2 - ^33_2^0$	O(0.15)1.44	
2490.40	$c^3P_2 - ^34_1^0$	O(0.00)1.50	
2467.06	$c^3P_2 - ^35_2^0$	O(0.00)1.49	
2459.46	$b^3P_1 - ^33_2^0$	O(0.36, 0.73)0.67, 1.04, - , 1.72	
2447.83	$b^3P_2 - ^33_2^0$	O(0.28)0.98	
2432.53	$a^3H_6 - x^3G_5^0$	O(0.00)1.09 C(0.00, 0.03, 0.07, 0.10, 0.14, 0.17) 1.00, 1.03, 1.07 1.33	Pattern may be unresolved
2429.93	$b^3P_1 - ^35_2^0$	O(0.00)1.53	
2418.50	$c^3P_1 - ^36_1^0$	O(0.00)1.04	
2416.22	$a^3H_5 - x^3G_4^0$	O(0.00)1.00 C(0.00, 0.03, 0.06, 0.09, 0.12) 0.90, 0.93, 0.96 1.14	Pattern may be unresolved
2386.74	$a^3G_5 - y^3F_4^0$	O(0.00)1.32 C(0.00, 0.05, 0.10, 0.15, 0.20) 1.00, 1.05, 1.10 1.40	Pattern may be unresolved
2360.51	$a^3G_4 - y^3F_3^0$	O(0.00)1.07 C(0.00 0.10)0.99 1.15	Pattern may be unresolved
2341.10	$a^3D_3 - ^38_3^0$	O(0.00, 1.10)	

Table XII. (Continued)

WAVE LENGTH	COMBINATION	ZEEMAN PATTERN	REMARKS
2337.00	$a^3G_5 - y^3G_5^0$	O(0.00)1.61 C(0.00)1.20	Discrepancy unexplained
2327.67	$a^3G_3 - y^3G_3^0$	O(0.00)1.00 C(0.00)0.75	Discrepancy unexplained
2256.73	$a^3H_6 - z^3H_6^0$	O(0.00)1.17 C(0.00)1.17	

From the observed Zeeman patterns, it appears that the line whose wave length is 2546.74 Å. does not arise from the combination $3d^7(^4P)4s\ b^3P_0 - 3d^7(^4P)4p\ z^5P_1^0$. The intensity of this line is anomalously large, and it must have fitted into the intersystem multiplet by coincidence.

The lines whose wave lengths are 2337.00 Å. and 2327.67 Å. exhibit Zeeman patterns that do not agree with their identifications as the combinations $3d^7(^2G)4s\ a^3G_5 - 3d^7(^2G)4p\ y^3G_5^0$, and $3d^7(^2G)4s\ a^3G_3 - 3d^7(^2G)4p\ y^3G_3^0$. Both of the 3G terms involved in these combinations are therefore to be considered questionable. This evidence is not, however, sufficient to warrant discarding either term, because it is only necessary for one of the terms to be incorrectly identified to produce this disagreement, and there is no way of telling which term may be faulty.

VI. ENERGY LEVELS

The known energy levels in the Co II spectrum are listed in Table XIII. They have been grouped together into terms whenever such grouping was justified, and the first two columns list the designations given to these terms and the configurations to which they have been assigned. There is a group of quintets which could not be assigned to a configuration on the basis of existing evidence. In cases where the levels could not be grouped into terms, they have been listed separately and designated by a sequence of numbers running from 2 to 17.

The observed g values for the new terms have been deduced from Weeks' data. Observed values for the previously established terms have been taken from Findlay's paper. The g values calculated from the Lande theory are presented in the table for comparison.

With the exception of the levels of the term $3d^8 a^3F$, all of those levels which are listed in the table with seven significant figures are consistent among themselves to within one-tenth of a wave number. The absolute values of these levels with respect to the normal state $3d^8 a^3F$ may be in error by several tenths of a wave number because all combinations with the normal state give rise to lines whose wave lengths fall in the vacuum region,

where measurements are less precise because the vacuum spectrograph used did not have as large dispersion as did the larger grating instrument used in the longer wave length region, and also because an error in wave length produces a larger error in wave number as the wave length decreases.

The following explanation pertains to the symbols used in Table XIII:

M.- Meggers

F - Findlay,

M*- Meggers, revised by Findlay

F*- Findlay, energy value adjusted
by Hager

H - Hager

Table XIII.
ENERGY LEVELS

CONFIG.	DESIG.	J	ENERGY LEVEL	ENERGY DIFF.	AUTHOR	g OBS.	g LANDE
$3d^8$	a^3F	4	0.00		F		
		3	950.28	950.28	F		
		2	1597.43	647.15	F		
$3d^7(^4F)4s$	a^5F	5	3349.86		M	1.39	1.40
		4	4028.39	678.53	M	1.36	1.35
		3	4560.23	531.84	M	1.25	1.25
		2	4949.48	389.25	M	1.00	1.00
		1	5204.09	254.61	M	0.00	0.00
$3d^7(^4F)4s$	b^3F	4	9812.24		F	1.28	1.25
		3	10707.82	895.58	F	1.10	1.08
		2	11321.35	613.53	F	0.60	0.67
$3d^8$	a^3P	2	13260.12		H	1.38	1.50
		1	13403.78	143.66	H	1.50	1.50
		0	13592.71	188.93	H	0.00	0.00
$3d^7(^4P)4s$	a^5P	3	17771.1		M*	1.70	1.67
		2	18031.0	259.9	M*	1.86	1.83
		1	18338.2	307.2	M*	2.48	2.50
$3d^7(^2G)4s$	a^3G	5	21624.2?		H		
		4	22008.9	384.7	H	1.06	1.05
		3	22226.2?	217.3	H		

Table XIII. (Continued)

CONFIG.	DESIG.	J	ENERGY LEVEL	ENERGY DIFF.	AUTHOR	$\frac{g}{\text{OBS.}}$	$\frac{g}{\text{LANDE}}$
$3d^7(^4P)4s$	b^3P	2	24074.2	193.1	H	1.49	1.50
		1	24267.3	143.8	H		
		0	24411.1		H	0.00	0.00
$3d^7(^2P)4s$	o^3P	2	24886.0	431.4	H		
		1	25317.4	543.8	H		
		0	25861.2		H		
$3d^7(^2H)4s$	a^3H	6	27105.2?	363.4	H	1.17	1.17
		5	27468.6?		H		
		4					
$3d^7(^2D)4s$	a^3D	3	27484.1	627.4	H	1.32	1.33
		2	28111.5	-526.5	H	1.17	1.17
		1	27585.0		H		
$3d^6 4s^2$	a^5D	4	40694.8	619.0	H		
		3	41313.8	424.0	H		
		2	41737.8	270.8	H		
		1	42008.6		H		
		0					

Table XIII. (Continued)

CONFIG.	DESIG.	J	ENERGY LEVEL	ENERGY DIFF.	AUTHOR	g OBS.	g LANDE
$3d^7(^4F)5s$	e^5F	5	84013.70		F*		
				572.65			
		4	84586.35		F*		
				580.75			
		3	85167.10		F*		
$3d^7(^4F)5s$	e^3F			428.60			
		2	85595.70		F*		
				280.14			
		1	85875.84		F*		
$3d^7(^4F)5s$	e^3F	4	85481.07		F*		
				865.12			
		3	86346.19		F*		
$3d^7(^4F)4p$	z^5F^0			593.41			
		2	86939.60		F*		
		5	45197.14		M	1.40	1.40
				181.06			
$3d^7(^4F)4p$	z^5F^0	4	45378.20		M*	1.38	1.35
				593.23			
		3	45971.43		M*	1.26	1.25
				480.76			
		2	46452.19		M*	1.02	1.00
$3d^7(^4F)4p$	z^5F^0			333.64			
		1	46785.83		M*	0.00	0.00
		4	46320.29		M*	1.48	1.50
				718.30			
$3d^7(^4F)4p$	z^5D^0	3	47038.59		M*	1.51	1.50
				498.36			
		2	47536.95		M*	1.52	1.50
				311.39			
		1	47848.34		M*	1.53	1.50
$3d^7(^4F)4p$	z^5D^0			146.77			
		0	47995.11		F	0.00	0.00

Table XIII. (Continued)

CONFIG.	DESIG.	J	ENERGY LEVEL	ENERGY DIFF.	AUTHOR	g OBS.	g LANDE
$3d^7(^4F)4p$	z^5G^0	6	47078.06		M	1.33	1.33
				267.29			
		5	47345.35		M	1.27	1.27
				461.66			
		4	47807.01		M	1.15	1.15
$3d^7(^4F)4p$	z^3G^0			343.40			
		3	48150.41		M	0.92	0.92
				237.43			
		2	48387.84		M	0.33	0.33
$3d^7(^4F)4p$	z^3G^0	5	48555.45		F	1.20	1.20
				792.35			
		4	49347.80		F	1.09	1.05
$3d^7(^4F)4p$	z^3F^0			687.94			
		3	50035.74		F	0.80	0.75
$3d^7(^4F)4p$	z^3F^0	4	49697.11		F	1.20	1.25
				684.01			
		3	50381.12		F	1.08	1.08
$3d^7(^4F)4p$	z^3D^0			532.73			
		2	50913.85		F	0.65	0.67
$3d^7(^4F)4p$	z^3D^0	3	51511.73		F	1.32	1.33
				717.49			
		2	52229.22		F	1.16	1.17
$3d^7(^4F)4p$	z^3D^0			454.88			
		1	52684.10		F	0.45	0.50
$3d^7(^4P)4p$	z^5S^0	2	56010.1		F	1.99	2.00

Table XIII. (Continued)

CONFIG.	DESIG.	J	ENERGY LEVEL	ENERGY DIFF.	AUTHOR	g OBS.	g LANDE
$3d^7(^4P)4p$	y^5D^0	4	61388.0		F	1.50	1.50
				-147.6			
		3	61240.4		M	1.51	1.50
				19.7			
		2	61260.1		M	1.51	1.50
				89.9			
		1	61350.0		F*		
				107.6			
		0	61457.6		F	0.00	0.00
$3d^7(^4P)4p$	z^3S^0	1	62439.6		H	2.03	2.00
$3d^7(^4P)4p$	z^5P^0	3	63344.0		M*	1.67	1.67
				23.0			
		2	63367.0		M*	1.86	1.83
				298.0			
		1	63665.0		M*	2.62	2.50
$3d^7(^4P)4p$	y^3D^0	3	63586.7		H		
				29.1			
		2	63615.8		H		
				249.2			
		1	63865.0		H		
$3d^7(^2G)4p$	y^3F^0	4	63510.2		H		
				850.0			
		3	64360.2		H		
				657.0			
		2	65017.2		H		
$3d^7(^2G)4p$	y^3G^0	5	64401.8?		H		
				752.4			
		4	65154.2?		H		
				20.6			
		3	65174.8?		H		

Table XIII. (Continued)

CONFIG.	DESIG.	J	ENERGY LEVEL	ENERGY DIFF.	AUTHOR	^g OBS.	^g LANDE
$3d^7(^2H)4p$	x^3G^o	5	68203.7?	640.4	H		
		4	68844.1?		H		
		3					
$3d^7(^2H)4p$	z^3H^o	6	71403.2?		H		
		5					
		4					
$3d^7_4p (?)$	$^3_2^o$	0	64605.3		H		
	$^3_3^o$	2	64914.3		H		
	$^3_4^o$	1	65028.2		H		
	$^3_5^o$	2	65408.0		H		
	$^3_6^o$	1	65657.0		H		
	$^3_7^o$	3	69060.0		H		
	$^3_8^o$	3	70186.0		H		
	$^3_9^o$	2	70265.8		H		
	$^3_{10}^o$	2	70457.5		H		
	$^3_{11}^o$	1	70774.9		H		
	$^3_{12}^o$	1	70856.8		H		
	$^3_{13}^o$	0	71281.1		H		
	$^3_{14}^o$	2	72653.4		H		
	$^3_{15}^o$	1	73112.4		H		
	$^3_{16}^o$	1	73143.9		H		
	$^3_{17}^o$	0	73823.2		H		

Table XIII. (Continued)

CONFIG.	DESIG.	J	ENERGY LEVEL	ENERGY DIFF.	AUTHOR	g OBS.	g LANDE
$3d^6 4s(^6D)4p$	x^5D^o	4	80298.8		H		
				243.7			
		3	80542.5		H		
				245.7			
		2	80788.2		H		
				182.7			
$3d^6 4s(^6D)4p$	y^5F^o	1	80970.9		H		
				95.4			
		0	81066.3?		H		
		5	81970.2		H		
				373.1			
$3d^6 4s(^6D)4p$	y^5F^o	4	82343.3		H		
				284.3			
		3	82627.6		H		
				332.0			
		2	82959.6		H		
		1					
$3d^6 4s(^6D)4p$	y^5P^o	3	85043.9		H		
				668.4			
		2	85712.3		H		
				421.7			
$3d^6 4s(^6D)4p$	y^5P^o	1	86134.0		H		
$3d^6 4s(^4D)4p$	$3D^o$	3	86981.0		H		
				508.7			
		2	87489.7		H		
		1					
$3d^6 4s(^4D)4p$	$3F^o$	4	87111.4?		H		
				784.7			
		3	87896.1		H		
				521.3			
$3d^6 4s(^4D)4p$	$3F^o$	2	88417.4		H		

Table XIII. (Continued)

CONFIG.	DESIG.	J	ENERGY LEVEL	ENERGY DIFF.	AUTHOR	^g OBS.	^g LANDE
$3d^6 4s(a^4D)4p$	$^3P^o$	2	90702.2		H		
		1	91289.0	586.8	H		
		0					
	$^5F^o$	5	99969.6		H		
				470.0			
		4	100439.6		H		
				578.7			
		3	101018.3		H		
				270.7			
		2	101289.0		H		
	$^5D^o$	4	101328.0		H		
				685.8			
		3	102013.8		H		
				413.5			
		2	102427.3		H		
		1					
		0					
	$^5D^o$	4	102264.7		H		
				493.6			
		3	102758.3		H		
				327.5			
		2	103085.8		H		
	$^5P^o$	3	98204.2		H		
				297.8			
		2	98502.0		H		
		1					

VII. IDENTIFIED LINES

Table XIV lists the wave length, estimated intensity, and wave number of each line that arises from combinations between established energy levels. These wave length measurements and intensity estimates are taken from the master list reproduced in Appendix III, which was compiled from many sets of measurements, and the method used for assembling the data for the master list is described in the appendix.

The following explanation pertains to the symbols used in Table XIV:

- F - Findlay
- M - Meggers
- M*- Meggers, revised by Findlay
- P - Combinations, arising from terms discovered by Meggers and Findlay, which appeared in Shenstone's vacuum measurements
- H - Hager
- w - Wide Line
- W - Very wide line
- CO I- Lines which may arise from the normal cobalt atom

Table XIV.
IDENTIFIED LINES

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
1010.971		0	98914.8	$a^5F_5 - v^5D_4^0$	H
1012.864		0	98729.9	$a^5F_4 - v^5D_3^0$	H
1265.930		4	78993.3	$a^5F_5 - y^5F_4^0$	H
71.936		5	620.3	$a^5F_5 - y^5F_5^0$	H
72.278		3	599.2	$a^5F_4 - y^5F_3^0$	H
76.901		4	314.6	$a^5F_4 - y^5F_4^0$	H
80.677		0	083.7	$b^3F_4 - 4s4p^3F_3^0$	H
80.950		1	067.1	$a^5F_3 - y^5F_3^0$	H
81.884		1	78010.2	$a^5F_2 - y^5F_2^0$	H
83.003		1	77942.1	$a^5F_4 - y^5F_5^0$	H
85.619		1	783.5	$a^5F_3 - y^5F_4^0$	H
86.083		0	755.5	$a^5F_1 - y^5F_2^0$	H
86.841		1	709.7	$b^3F_3 - 4s4p^3F_2^0$	H
87.361		1	678.3	$a^5F_2 - y^5F_3^0$	H
93.671		6	299.4	$b^3F_4 - 4s4p^3F_4^0$	H
95.539		4	188.0	$b^3F_3 - 4s4p^3F_3^0$	H
95.864		7	168.6	$b^3F_4 - 4s4p^3D_3^0$	H
97.094		3	77095.4	$b^3F_2 - 4s4p^3F_2^0$	H
1299.574		8	76948.3	$a^5F_5 - x^5D_4^0$	H
1302.390		6	76781.9	$b^3F_3 - 4s4p^3D_2^0$	H

Table XIV. (Continued)

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
1306.966		8	76513.1	$a^5F_4 - x^5D_3^0$	H
11.117		5	270.8	$a^5F_4 - x^5D_4^0$	H
11.856		5	227.9	$a^5F_3 - x^5D_2^0$	H
12.869		2	169.1	$b^3F_2 - 4s4p^3D_2^0$	H
15.419		4	76021.4	$a^5F_2 - x^5D_1^0$	H
16.077		5	75983.0	$a^5F_3 - x^5D_3^0$	H
18.180		2	862.2	$a^5F_1 - x^5D_0^0$	H
18.586		3	838.8	$a^5F_2 - x^5D_2^0$	H
1319.836		3	75767.0	$a^5F_1 - x^5D_1^0$	H
1424.784	0	5	70186.1	$a^3F_4 - ^38_3^0$	H
38.693		4	69507.5	$a^3F_3 - ^310_2^0$	H
44.335		2	236.0	$a^3F_3 - ^38_3^0$	H
45.544		3	178.1	$a^3F_2 - ^311_1^0$	H
48.009	2	10	69060.3	$a^3F_4 - ^37_3^0$	H
52.558		1	68844.1	$a^3F_4 - x^3G_4^0$	H
66.196	3	10	203.7	$a^3F_4 - x^3G_5^0$	H
68.204		6	110.4	$a^3F_3 - ^37_3^0$	H
68.356	1w	6	68103.4	$a^5P_2 - y^5P_1^0$	H
71.860	1	6	67941.2	$a^5P_3 - y^5P_2^0$	H
72.892	4	8	894.0	$a^3F_3 - x^3G_4^0$	H
1475.018		2	67795.8	$a^5P_1 - y^5P_1^0$	H

Table XIV. (Continued)

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
1484.251		5	67374.1	$a^5P_1 - y^5P_2^0$	H
86.483	4	9	272.9	$a^5P_3 - y^5P_3^0$	H
87.768	2	2	214.8	$b^3P_2 - 4s4p^3P_1^0$	H
92.064		1	021.3	$b^3P_1 - 4s4p^3P_1^0$	H
92.254	0	6	67012.3	$a^5P_2 - y^5P_3^0$	H
1495.256		2	66878.2	$b^3P_0 - 4s4p^3P_1^0$	H
1500.875	1	5	627.8	$b^3P_2 - 4s4p^3P_2^0$	H
05.234		2	66434.9	$b^3P_1 - 4s4p^3P_2^0$	H
19.391		4	65815.8	$c^3P_2 - 4s4p^3P_2^0$	H
34.820		3	65154.2	$a^3F_4 - y^3G_4^0$	H
52.752		9	64401.8	$a^3F_4 - y^3G_5^0$	H
53.754		4	360.3	$a^3F_4 - y^3F_3^0$	H
57.039		6	224.5	$a^3F_3 - y^3G_3^0$	H
57.542		6	203.6	$a^3F_3 - y^3G_4^0$	H
60.865		4	64067.0	$a^3F_3 - y^3F_2^0$	H
72.645	1	7	63587.1	$a^3F_4 - y^3D_3^0$	H
72.886	0	4	577.4	$a^3F_2 - y^3G_3^0$	H
74.548	0	8	510.3	$a^3F_4 - y^3F_4^0$	H
76.789	2	7	420.0	$a^3F_2 - y^3F_2^0$	H
77.042	2	7	409.9	$a^3F_3 - y^3F_3^0$	H
1581.715		1	63222.5	$b^3P_1 - 4s4p^3D_2^0$	H

Table XIV. (Continued)

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
1589.654		3	62906.8	$b^3P_2 - 4s4p^3D_3^0$	H
93.299		5	762.9	$a^3F_2 - y^3F_3^0$	H
93.439		2	757.3	$a^5P_2 - x^5D_2^0$	H
95.773	0	6	665.6	$a^3F_3 - y^3D_2^0$	H
96.508		3	636.7	$a^3F_3 - y^3D_3^0$	H
96.614		1	632.6	$a^5P_1 - x^5D_1^0$	H
98.470		5	559.8	$a^3F_3 - y^3F_4^0$	H
99.308		4	527.0	$a^5P_3 - x^5D_4^0$	H
1599.701		3	511.7	$a^5P_2 - x^5D_3^0$	H
1601.282		2	450.0	$a^5P_1 - x^5D_2^0$	H
05.962		5	268.0	$a^3F_2 - y^3D_1^0$	H
12.410		2	62019.0	$a^3F_2 - y^3D_2^0$	H
24.024	0	4	61570.0	$a^5D_4 - v^5D_4^0$	H
27.482		3	444.6	$a^5D_3 - v^5D_3^0$	H
30.044		1	348.0	$a^5D_2 - v^5D_2^0$	H
37.269		3	077.3	$a^5D_1 - v^5D_2^0$	H
38.787	0w	4	61020.7	$a^5D_2 - v^5D_3^0$	H
40.669		4	60950.8	$a^5D_3 - v^5D_4^0$	H
47.446	0	6	700.0	$a^5D_3 - w^5D_3^0$	H
47.729		3	689.6	$a^5D_2 - w^5D_2^0$	H
1649.263	1	7	60633.1	$a^5D_4 - w^5D_4^0$	H

Table XIV. (Continued)

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
1655.119		3	60418.6	$a^5D_1 - w^5D_2^0$ $a^3P_2 - ^317_0^0$	H H
59.032		4	276.1	$a^5D_2 - w^5D_3^0$	H
66.276		5	60014.1	$a^5D_3 - w^5D_4^0$	H
70.799	0	2	851.6	$a^3P_3 - ^316_1^0$	H
73.787		1	744.8	$a^5D_4 - x^5F_4^0$	H
73.918		3	740.1	$a^3P_1 - ^316_1^0$	H
74.921	1	5	704.3	$a^5D_3 - x^5F_3^0$	H
77.093	0	1	627.0	$a^3D_3 - 4s4p^3F_4^0$	H
79.232		2	60551.0	$a^5D_2 - x^5F_2^0$ $a^3P_0 - ^316_0^0$	H H
80.125		3	59519.4	$a^3P_1 - ^315_1^0$	H
83.722	3	6	392.2	$a^3P_2 - ^314_2^0$	H
86.896	1	5	280.5	$a^5D_2 - x^5F_2^0$ $a^5D_1 - x^5F_3^0$	H H
87.057	1	5	274.8	$a^5D_4 - x^5F_5^0$	H
87.788		1	249.1	$a^3P_1 - ^314_2^0$	H
91.310	2	4	125.8	$a^5D_3 - x^5F_4^0$	H
1693.330	1	3	59055.2	$a^5F_3 - y^3D_2^0$	H
1700.487		1	58806.7	$a^5F_3 - z^5P_2^0$	P
12.614		1	390.3	$b^3F_4 - x^3G_5^0$	H
23.013		4	58037.9	$a^5F_5 - y^5D_4^0$	P
27.806	1	3	57876.9	$a^3P_1 - ^313_0^0$	H
1736.225		2	57596.2	$a^3P_2 - ^312_1^0$	H

Table XIV. (Continued)

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
1743.381	0	3	57359.8	$a^5F_4 - y^5D_4^0$	P
46.297		2	264.0	$a^3P_0 - ^312_1^0$	H
47.882		5	212.0	$a^5F_4 - y^5D_3^0$	P
48.343		3	197.0	$a^3P_2 - ^310_2^0$	H
54.217		7	57005.5	$a^3P_2 - ^39_2^0$	H
58.642		6	56862.1	$a^3P_1 - ^39_2^0$	H
63.675		5	699.8	$a^5F_3 - y^5D_2^0$	P
64.279		4	680.4	$a^5F_3 - y^5D_3^0$	P
73.040	1	4	400.3	$a^5F_2 - y^5D_1^0$	P
75.867		3	310.5	$a^5F_2 - y^5D_2^0$	P
77.658		3	253.8	$a^5F_1 - y^5D_0^0$	P
81.068	1	2	146.1	$a^5F_1 - y^5D_1^0$	P
1783.909		2	56056.7	$a^5F_1 - y^5D_2^0$	P
1833.271		1	54547.3	$b^3F_4 - y^3F_3^0$	H
55.037	2	2 CO I	53907.3	$b^3F_3 - y^3D_2^0$	H
59.625		2	774.3	$b^3F_4 - y^3D_3^0$	H
1862.353		6	53695.5	$b^3F_2 - y^3F_2^0$	H
1908.518	0	6	52396.7	$a^3P_2 - ^36_1^0$	H
13.756		2	253.3	$a^3P_1 - ^36_1^0$	H
17.627	1	7	147.8	$a^3P_2 - ^35_2^0$	H
1920.706		3	52064.2	$a^3P_0 - ^36_1^0$	H

Table XIV. (Continued)

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
1922.927		5	52004.1	$a^3P_1 - {}^35_2^0$	H
35.963		1	51653.9	$a^3P_2 - {}^33_2^0$	H
37.069		4	624.4	$a^3P_1 - {}^34_1^0$	H
40.64		50 F	512.3	$a^3F_4 - z^3D_3^0$	F
44.188		3	435.4	$a^3P_0 - {}^34_1^0$	H
50.098		15	279.5	$a^3F_3 - z^3D_2^0$	F
53.071		3	201.4	$a^3P_1 - {}^32_0^0$	H
57.424		14	51087.6	$a^3F_2 - z^3D_1^0$	F
75.008		4	50632.7	$a^3F_2 - z^3D_2^0$	F
81.736		2	460.8	$a^3P_1 - y^3D_1^0$	H
83.953		4	404.4	$a^3P_2 - z^5P_1^0$	H
84.843		3	381.8	$a^3F_4 - z^3F_3^0$	F
87.030		4	326.4	$a^3P_2 - y^3D_3^0$	H
91.567		3	211.7	$a^3P_1 - y^3D_2^0$	H
95.733		2	106.9	$a^3P_2 - z^5P_2^0$	H
97.135		3	071.7	$a^3P_0 - z^5P_1^0$	H
1998.549		1	50036.3	$a^3F_4 - z^5G_3^0$	F
2001.445	κ	7	49963.9	$a^3F_3 - z^3F_2^0$	F
08.877		4	779.1	$a^3G_5 - z^3H_6^0$	H
12.199	0	16	697.18	$a^3F_3 - z^3F_4^0$	F
2023.019	0	15	49430.84	$a^3F_3 - z^3F_3^0$	M*

Table XIV. (continued)

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2026.424	0	13	49347.73	$a^3F_4 - z^3G_4^0$	M*
27.703	0	14	316.42	$a^3F_2 - z^3F_2^0$	M*
33.379		10	179.2	$a^3P_2 - z^3S_1^0$	H
37.244		11	085.9	$a^3F_3 - z^3G_3^0$	M*
39.333		9	49035.6	$a^3P_1 - z^3S_1^0$	H
47.223		3	48846.7	$a^3P_0 - z^3S_1^0$	H
49.831		5	784.5	$a^3F_2 - z^3F_3^0$	M*
51.42		10	746.8	$a^3F_3 - z^3F_4^0$	F
51.999		3	733.0	$b^3P_0 - ^316_1^0$	H
59.480		11	555.9	$a^3F_4 - z^3G_5^0$	M*
64.453		10	439.0	$a^3F_2 - z^3G_3^0$	M*
72.209		3	257.7	$c^3P_2^0 - ^316_1^0$	H
79.485		1	48088.8	$a^3P_2 - y^5D_1^0$	H
83.358	Ow	10	47999.70	$a^3P_2 - y^5D_2^0$	H
90.889		1	826.5	$c^3P_1 - ^316_1^0$	H
91.738		4	807.1	$a^3F_4 - z^5G_4^0$	M*
94.931		2	734.3	$a^5F_2 - z^3D_1^0$	M*
97.809		3	668.8	$a^5F_3 - z^3D_2^0$	M*
2099.230		2	636.5	$a^5P_3 - ^35_2^0$	H
2106.019		3	47483.0	$a^5F_4 - z^3D_3^0$	M*

Table XIV. (Continued)

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2106.162		3	47479.7	$a^5F_1 - z^3D_1^0$	M*
12.140		7	345.3	$a^3F_4 - z^5G_4^0$	M*
12.540		1	336.4	$o^3P_1 - {}^3\lambda_2^0$	H
13.335		2	318.1	$a^5P_1 - {}^3\epsilon_1^0$	H
14.966	0	4	282.10	$c^3P_0 - {}^3\epsilon_1^0$	H
15.088		2	279.4	$a^5F_2 - z^3D_2^0$	M*
18.644		4	200.0	$a^3F_3 - z^5G_3^0$	M*
24.524		2	069.4	$a^5P_1 - {}^3\epsilon_2^0$	H
26.532		3	47024.9	$a^5F_1 - z^3D_2^0$	M*
29.87		1	46951.3	$a^5F_3 - z^3D_3^0$	M*
34.174		4	856.5	$a^3F_3 - z^5G_4^0$	M*
37.178		3	790.7	$a^3F_2 - z^5G_2^0$	M*
46.932		6	561.8	$a^5F_2 - z^3D_3^0$	M*
48.06		2	553.6	$a^3F_2 - z^5G_3^0$	M*
57.381		10 M	352.5	$a^5F_3 - z^3F_2^0$	M*
				$a^5F_4 - z^3F_3^0$	M*
57.635		40 M	347.0	$a^5F_5 - z^3F_4^0$	M*
58.862		1	320.7	$a^3F_4 - z^5D_4^0$	M*
62.67		2 M	239.1	$a^3D_1 - {}^3\lambda_0^0$	H
69.766	0	3	087.90	$a^3F_3 - z^5D_3^0$	H
73.573		3	46007.2	$a^5F_4 - z^3G_3^0$	M*
2174.008	1	9	45998.40	$a^5F_5 - z^3G_4^0$ $b^3P_1 - {}^3\epsilon_2^0$	M* H

Table XIV. (Continued)

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2175.614		2	45964.0.	$a^5F_2 - z^3F_2^0$ $o^3P_2 - z^3D_2^0$	M* H
81.299		6	844.3	$a^5P_3 - y^3D_2^0$	H
82.414		4	820.8	$a^5F_3 - z^3F_3^0$	F
82.679		7	815.3	$a^5P_3 - y^3D_3^0$	H
87.730		25 M	709.5	$a^5F_1 - z^3F_2^0$	M*
89.686		25 M	668.6	$a^5F_4 - z^3F_4^0$	M*
91.34		75 M	634.2	$a^5P_2 - z^5P_1$	M*
91.638		7	628.0	$a^3D_3 - z^3D_1^0$	H
93.180		11	595.9	$a^5P_3 - z^5P_2^0$	M*
94.293	3w	12	572.85	$a^5P_3 - z^5P_3^0$	M*
95.89		1 M	539.7	$o^3P_1 - z^3D_1^0$	H
96.515	0	4	526.56	$a^5P_1 - y^3D_1^0$ $a^3D_1 - z^3D_1^0$	H H
2198.968		6	475.9	$a^5F_3 - z^3G_3^0$	M*
2201.101	0	25 M	431.79	$a^5F_2 - z^3F_3^0$	M*
03.617		11	379.9	$o^3P_2 - z^3G_2^0$ $a^3F_4 - z^5F_4^0$	H M*
04.106		3	369.9	$a^3F_3 - z^5D_4^0$	M*
05.750		8	336.1	$a^5P_2 - z^5P_2^0$	H
06.205		9	326.7	$a^5P_1 - z^5P_1^0$	M*
06.576		9	319.1	$a^5F_4 - z^3G_4^0$	M*
06.905	4w	75 M	312.99	$a^5P_2 - z^5P_3^0$	M*
2212.102		10	45205.9	$a^5F_5 - z^3G_5^0$	M*

Table XIV. (continued)

WAVE LENGTH IN VACUUM	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2212.902	0	2	45189.27	$a^3F_2 - z^5F_1^0$	H
13.825		7	170.1	$a^3D_3 - ^314_2^0$	H
15.456	1	20 M	137.11	$a^5F_3 - z^3F_4^0$	M*
17.968	3w	10	086.10	$a^5F_2 - z^3G_3^0$	M*
20.594	1	8	032.41	$a^3D_2 - ^316_1^0$	H
20.775	2	10	028.89	$a^5P_1 - z^5P_2^0$	M*
21.171	0w	4	020.83	$a^3F_3 - z^5F_3^0$	H
22.123	1	9	45001.44	$a^3D_2 - ^315_1^0$	H
22.422	2	9	44995.54	$a^3P_0 - ^312_1^0$	H
22.915		4	986.0	$b^3P_2 - ^37_3^0$	H
2224.745	1	6	44948.49	$a^3P_1 - ^39_2^0$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2232.065	3	50	44787.60	$a^5F_3 - z^3G_4^0$	M*
45.120	15	100	527.20	$a^5F_4 - z^3G_5^0$	M*
48.66		8	457.1	$a^5F_5 - z^5G_4^0$	M*
51.10		10	408.9	$a^5P_2 - z^3S_1^0$	H
51.63		2	398.5	$a^5D_3 - y^5P_2^0$	H
54.181	0	1	348.13	$a^5D_4 - y^5P_3^0$	H
65.76		6	121.7	$a^5F_4 - z^5G_3^0$	M*
66.795	1w	10	44101.47	$a^5P_1 - z^3S_1^0$	H
72.265	0	20	43995.42	$a^5F_5 - z^5G_5^0$	M*
80.938	0	7	828.05	$a^5F_3 - z^5G_2^0$	M*
83.511	3w	20	778.77	$a^5F_4 - z^5G_4^0$	M*
86.147	50	150	728.20	$a^5F_5 - z^5G_6^0$	M*
87.82		4 CO I	696.2	$a^3D_1 - ^313_0^0$	H
91.981	12	40w	616.91	$a^5P_3 - y^5D_4^0$	M*
93.411	10w	30	589.71	$a^5F_3 - z^5G_3^0$	M*
98.727	0	10	488.92	$a^5P_3 - y^5D_2^0$	M*
2299.76		25	469.4	$a^5P_3 - y^5D_3^0$	M*
2301.401	2	15	438.39	$a^5F_2 - z^5G_2^0$	M*
07.507	0	2	323.46	$a^5P_2 - y^5D_1^0$	M*
07.854	25	75	317.04	$a^5F_4 - z^5G_5^0$	M*
2311.596	25	50	43246.81	$a^5F_3 - z^5G_4^0$	M*

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2312.548	6	10	43229.02	$a^5P_2 - y^5D_1^0$	M*
13.619	3w	8	209.02	$a^5P_2 - y^5D_3^0$	M*
14.053	25	40	200.91	$a^5F_2 - z^5G_3^0$	M*
14.646	1	1	189.84	$a^3D_1 - ^311_1^0$	H
14.969	40	30	183.81	$a^5F_1 - z^5G_2^0$	M*
18.420	2w	8	119.21	$a^5P_1 - y^5D_0^0$	M*
24.310	10	40	43010.29	$a^5P_1 - y^5D_1^0$	M*
				$a^5F_4 - z^5D_3^0$	M*
26.103	35	20	42977.14	$a^5F_3 - z^5D_2^0$	M*
26.458	30	25	970.59	$a^5F_5 - z^5D_4^0$	M*
29.094	8w	10	921.96	$a^5P_1 - y^5D_2^0$	M*
30.361	15	30	898.81	$a^5F_2 - z^5D_1^0$	M*
36.223	15	20	790.99	$a^5F_1 - z^5D_0^0$	M*
				$a^3G_3 - y^5F_2^0$	H
38.72		6	745.2	$a^3D_2 - ^312_1^0$	H
41.112	15	40	701.64	$a^3D_3 - ^38_3^0$	H
42.247	0w	2	680.95	$a^3D_1 - ^39_2^0$	H
44.254	15	25	644.42	$a^5F_1 - z^5D_1^0$	M*
47.382	25	30	587.60	$a^5F_2 - z^5D_2^0$	M*
53.386	75	60	478.95	$a^5F_3 - z^5D_3^0$	M*
60.508	25	20	350.80	$a^3G_4 - y^3F_3^0$	H
60.784	20	8	345.84	$a^3D_2 - ^310_2^0$	H
2361.510	8	10	42332.83	$a^5F_1 - z^5D_2^0$	M*

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2363.802	100	90	42291.78	$a^5F_4 - z^5D_4^0$	M
75.190	15	15	089.05	$a^5F_2 - z^5D_3^0$	M
76.023	0		074.28	$a^3D_2 - {}^3P_3^0$	H
78.620	100	100	42028.35	$a^5F_5 - z^5F_4^0$	M
83.455	100	80	41943.10	$a^5F_4 - z^5F_3^0$	M
86.367	75	50	891.92	$a^5F_3 - z^5F_2^0$	M
86.717	75w	25	885.83	$a^3G_5 - y^3F_4^0$	H
88.918	100	100	847.19	$a^5F_5 - z^5F_5^0$	M
89.537	30	40	836.35	$a^5F_2 - z^5F_1^0$	M
93.901	10	10	760.09	$a^5F_3 - z^5D_4^0$	M
2397.380	20	60	699.50	$b^3F_4 - z^3D_3^0$	M
2404.106	25	20	581.65	$a^5F_1 - z^5F_1^0$	M
04.523	10w	10	575.63	$a^3D_3 - {}^3P_3^0$	H
07.664	10	20	521.39	$b^3F_3 - z^3D_2^0$	M
08.746	15	25	502.74	$a^5F_2 - z^5F_2^0$	M
08.833	3	0?	501.24	$a^3G_4 - y^3F_4^0$	H
14.068	50	40	411.26	$a^5F_3 - z^5F_3^0$	M
16.214	25	15	374.48	$a^3H_5 - x^3G_4^0$	H
16.899	30	30	362.74	$b^3F_2 - z^3D_1^0$	M
17.659	75	40	349.75	$a^5F_4 - z^5F_4^0$	M
2418.584	1		41333.93	$b^3P_2 - {}^3P_2^0$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2422.04		1	41274.8	$a^5D_4 - y^5F_5^0$	H
23.626	10	10	247.96	$a^5F_1 - z^5F_2^0$	M
23.756	2	3	245.75	$b^3P_0 - {}^36_1^0$	H
28.289	10	10	168.77	$a^5F_4 - z^5F_5^0$	M
29.928	4	5	140.99	$b^3P_1 - {}^35_2^0$	H
32.517	35	30	097.20	$a^3H_6 - x^3G_5^0$	H
36.983	4	10	41021.90	$a^5F_2 - z^5F_3^0$	M
43.781	15	40	40907.81	$b^3F_2 - z^3D_2^0$	M
47.826	8	12	840.20	$b^3P_2 - {}^33_2^0$	H
49.166	6	10	817.86	$a^5F_3 - z^5F_4^0$	M
50.005	12	35	803.88	$b^3F_3 - z^3D_3^0$	M
52.556	1w	2	761.44	$b^3P_1 - {}^34_1^0$	H
54.14		1	735.30	$a^3H_5 - x^3G_5^0$	H
59.452	10	10	647.16	$b^3P_1 - {}^33_2^0$	H
64.200	20	35	568.85	$b^3F_4 - z^3F_3^0$	M
67.051	15	20	521.97	$c^3P_2 - {}^35_2^0$	H
78.195	6	4	339.75	$c^3P_1 - {}^36_1^0$	H
78.299	5	5	338.06	$b^3P_1 - {}^32_0^0$	H
81.534	2	0	285.49	$b^3P_2 - y^3F_3^0$	H
85.358	8	10	223.51	$b^3F_3 - z^3G_3^0$	H
2486.439	15	35	40206.02	$b^3F_4 - z^3F_2^0$	M

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2487.150	3	5	40194.52	$b^3P_0 - {}^3_2D_0$	H
87.376	3w	4	190.88	$b^3F_2 - z^3D_3$	M
90.392	10w	15	142.16	$c^3P_2 - {}^3_4D_1$	H
93.572	0	0	091.03	$c^3P_1 - {}^3_5D_2$	H
2497.490	6w	10	40028.14	$c^3P_2 - {}^3_3D_2$	H
2506.473	75	70	39884.68	$b^3F_4 - z^3F_4^0$	M
12.061	5w	3	796.00	$c^3P_0 - {}^3_6D_1$	H
12.406	2w	2	790.54	$b^3P_2 - y^3D_1^0$	H
17.439	4	3	710.95	$c^3P_1 - {}^3_4D_1$	H
19.823	40	60	673.39	$b^3F_3 - z^3F_3^0$	M
24.212	1		604.45	$a^5D_4 - x^5D_4^0$	H
24.643	8w	7	597.68	$b^3P_1 - y^3D_1^0$	H
				$c^3P_1 - {}^3_3D_2^0$	H
24.971	75	80	592.55	$b^3F_2 - z^3F_2^0$	M
28.216	8w	5	541.70	$b^3P_2 - y^3D_2^0$	H
28.613	75	50	535.49	$b^3F_4 - z^3G_4^0$	M
30.105	50w	40	512.18	$b^3P_2 - y^3D_3^0$	H
33.821	35w	30	454.23	$b^3P_0 - y^3D_1^0$	H
37.454	1	3	397.75	$b^3P_1 - z^5P_1^0$	H
40.621	25	20	348.86	$b^3P_1 - y^3D_2^0$	H
41.958	30	50	327.95	$b^3F_3 - z^3G_3^0$	M
2544.267	75 CO I		39292.26	$b^3P_2 - z^5P_2^0$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2544.541	2w	4	39288.03	$c^3P_1 - z^3D_0^0$	H
45.040	2w	5	280.33	$a^3P_1 - z^3D_1^0$	H
45.718	1	2	269.8	$b^3P_2 - z^5P_3^0$	H
46.743	20	20	254.06	$b^3P_0 - z^5P_1^0$	H
49.723	4	OW	208.19	$z^5F_4^0 - e^5F_5$	F
50.534	2		195.72	$z^5F_3^0 - e^5F_3$	F
52.391	4	4	167.20	$c^3P_0 - z^3D_1^0$	H
53.938	3	OW	143.48	$z^5F_2^0 - e^5F_2$	F
56.761	25	8 CO I	100.26	$b^3P_1 - z^5P_2^0$	H
57.341	6	8	091.39	$a^3P_0 - z^3D_1^0$	H
57.438	1		089.91	$z^5F_1^0 - e^5F_1$	H
59.410	60	40	39059.80	$b^3F_2 - z^3F_3^0$	M
64.044	100	75	38989.21	$b^3F_3 - z^3F_4^0$	M
64.293	0		985.42	$a^5D_3 - x^5D_4^0$	H
65.367	2	6	969.10	$a^3P_2 - z^3D_2^0$	H
74.860	8	10	825.44	$a^3P_1 - z^3D_2^0$	H
75.474	20	OW	816.19	$z^5F_5^0 - e^5F_5$	F
75.893	1		809.87	$z^5F_2^0 - e^5F_1$	H
76.206	3		805.16	$a^5D_2 - x^5D_3^0$	H
77.927		1	779.3	$c^3P_2 - z^5P_1^0$	H
2580.335	50	100	38743.07	$b^3F_4 - z^3G_5^0$	M

Table XIV. (continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2581.211	2w	2w	38729.92	$c^3P_2 - y^3D_2^0$	H
82.244	30	50	714.42	$b^3F_2 - z^3G_3^0$	M
83.170	15	15	700.55	$c^3P_2 - y^3D_3^0$	H
87.223	50	60	639.93	$b^3F_3 - z^3G_4^0$	M
87.521	12		635.48	$z^5F_4^0 - e^5F_5$	F
88.900	10	0w	614.90	$z^5F_3^0 - e^5F_4$	F
92.709	1	2w	558.18	$z^5D_3^0 - e^5F_2$	F
2593.429	0w	2	547.47	$c^3P_1 - y^3D_1^0$	H
2605.724		20	365.7	$b^3P_2 - z^3S_1^0$	H
06.950		1	347.6	$c^3P_1 - z^5P_3^0$	H
07.600		0	338.0	$b^3F_4 - z^5G_3^0$	H
10.300		1	298.4	$c^3P_1 - y^3D_2^0$	H
12.499	15	1w	266.11	$z^5D_4^0 - e^5F_4$	F
13.494	25	10	251.54	$a^3P_2 - z^3D_3^0$	H
14.353	30w	20	238.97	$a^5P_3 - z^5S_2^0$	M
18.894	8	20	172.4	$b^3P_1 - z^3S_1^0$	H
21.924	8		128.57	$z^5D_3^0 - e^5F_3$	F
26.728	8	2w	058.83	$z^5D_2^0 - e^5F_2$	F
28.83		6	028.4	$b^3P_0 - z^3S_1^0$	H
28.886	2		027.60	$z^5D_1^0 - e^5F_1$	F
2630.540		3	38003.6	$c^3P_0 - y^3D_1^0$	H

TABLE XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2632.239	20w	30	37979.16	$a^5P_2 - z^5S_2^0$	M
36.070	12w	6	923.96	$a^3D_3 - {}^3S_2^0$	H
39.081	3	1W	880.69	$z^5D_0^0 - e^5F_1$	F
43.108		2	823.0	$a^3D_1 - {}^3S_2^0$	H
44.464	0w	1	803.59	$c^3P_0 - z^5P_1^0$	H
48.407	3		747.32	$z^5D_1^0 - e^5F_2$	H
52.193	8	5W	693.44	$z^5D_4^0 - e^5F_5$	F
53.095	0	1	680.63	$b^3F_3 - z^5G_2^0$	M
53.713	12w	15	671.85	$a^5P_1 - z^5S_2^0$	M
55.05		0	652.9	$a^3P_2 - z^3F_2^0$	H
56.641	3w		630.33	$z^5D_2^0 - e^5F_3$	H
62.056		1	553.8	$c^3P_2 - z^3S_1^0$	H
62.482	4		547.78	$z^5D_3^0 - e^5F_4$	F
62.640	0	1	545.55	$a^3D_2 - {}^3G_1^0$	H
63.534	30	60	532.95	$b^3F_4 - z^5G_5^0$	M
65.344	0		507.46	$a^3P_1 - z^3F_2^0$	H
66.728	5	5w	488.00	$z^5G_2^0 - e^5F_1$	F
69.773	20	4 W	445.24	$z^5G_3^0 - e^5F_2$	F
70.831	2W	2w	430.42	$a^3D_3 - {}^3S_2^0$	H
75.856	20	15W	360.13	$z^5G_4^0 - e^5F_3$	H
2680.435	0	1	37296.35	$a^3D_2 - {}^3S_2^0$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2681.903	0	1w	37275.90	$b^3P_2 - z^5D_1^0$	H
84.421	20	15W	240.94	$z^5G_5^0 - e^5F_4$	F
88.41		0	185.7	$b^3P_2 - y^5D_2^0$	H
89.826		2	166.1	$b^3P_2 - y^5D_3^0$	H
92.982	1	1	122.55	$o^3P_1 - z^3S_1^0$	H
93.090	1	2	121.06	$a^3P_2 - z^3F_3^0$	H
94.680	25	30	099.16	$b^3F_3 - z^5G_4^0$	M
95.857		5	082.9	$b^3P_1 - y^5D_1^0$	H
2697.069		8	066.3	$b^3F_2 - z^5G_2^0$	M
2700.415	0	2	020.38	$a^3D_1 - ^32^0$	H
00.686	1		37016.66	$z^5G_3^0 - e^5F_3$	H
02.019	8	2W	36998.41	$z^3G_4^0 - e^3F_3$	F
02.424	0		992.86	$b^3P_1 - y^5D_2^0$	H
06.507	30	10W	937.06	$z^5G_6^0 - e^5F_5$ $b^3P_0 - y^5D_1^0$	F H
07.348	20	8W	925.58	$z^3G_5^0 - e^3F_4$	F
08.002	2	8	916.67	$a^3D_2 - ^34_1^0$	H
08.943	8	1W	903.84	$z^3G_3^0 - e^3F_2$	F
14.437	6w	20	829.15	$b^3F_3 - z^5D_2^0$ $b^3F_2 - z^5G_3^0$	M M
16.380	0	5	802.80	$a^3D_2 - ^33_2^0$	H
18.111	1		779.37	$z^5G_4^0 - e^5F_4$	H
2718.423		0	36775.2	$a^3P_2 - z^3G_3^0$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2727.773	10	10W	36649.10	$z^3F_4^0 - e^3F_3$	F
33.009	0	6	578.90	$c^3P_0 - z^3S_1^0$	H
34.539	8	5W	558.44	$z^3F_3^0 - e^3F_2$	F
36.91		2	526.8	$b^3F_2 - z^5D_1^0$	M
38.322		4	507.80	$b^3F_4 - z^5D_4^0$	F
41.656		0	463.5	$c^3P_2 - y^5D_1^0$	H
48.433		1	373.6	$c^3P_2 - y^5D_2^0$	H
49.909		3	354.1	$c^3P_2 - y^5D_3^0$	H
51.688		0	330.6	$b^3F_3 - z^5D_3^0$	H
53.216	6	2W	310.45	$z^3G_3^0 - e^3F_3$	F
55.511	0	0	280.21	$a^3D_1 - y^3D_1^0$	H
60.443	0	0	215.39	$b^3F_2 - z^5D_2^0$	H
64.721	1	0	159.35	$b^3F_4 - z^5F_3^0$	F
66.719	15	8W	133.24	$z^3G_4^0 - e^3F_4$	F
69.072	1w	10	102.54	$a^3D_3 - y^3D_3^0$	H
70.818		0	079.8	$a^3D_1 - z^5P_1^0$	H
74.578	0	0	030.90	$a^3D_1 - y^3D_2^0$	H
74.964	15	2W	36025.89	$z^3F_2^0 - e^3F_2$	F
79.660	5	5W	35965.03	$z^3F_3^0 - e^3F_3$	F
81.399		0	942.5	$c^3P_1 - y^5D_2^0$	H
2786.049		3	35882.6	$a^3D_3 - z^5P_2^0$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2787.832		1	35859.6	$a^3D_3 - z^5P_3^0$	H
93.726	15	5W	783.99	$z^3F_4^0 - e^3F_4$	F
96.134		2	753.1	$a^3D_2 - y^3D_1^0$	H
96.837		3	744.2	$b^3F_3 - z^5F_2^0$	F
2798.949	0w	3w	717.19	$b^3F_2 - z^5D_3^0$	F
2807.173	1	10	612.55	$b^3F_3 - z^5D_4^0$	M
10.853	8	12	565.93	$b^3F_4 - z^5F_4^0$	M
15.745	0	3	504.14	$a^3D_2 - y^3D_2^0$	H
16.971		0	488.7	$c^3P_0 - y^5D_1^0$	H
18.074		0	474.8	$a^3D_2 - y^3D_3^0$	H
18.900	1w	5	464.41	$b^3F_2 - z^5F_1^0$	M
19.382	0w		458.34	$z^3G_5^0 - e^5F_5$	H
20.402	0		445.53	$z^3G_3^0 - e^3F_4$	H
21.449	5	1W	432.37	$z^3F_2^0 - e^3F_3$	F
25.235	2	15	384.89	$b^3F_4 - z^5F_5^0$	M
34.942	2	10	263.74	$b^3F_3 - z^5F_3^0$	M
37.480		0	232.2	$a^3D_2 - z^5P_3^0$	H
45.688	"	5	130.6	$a^3F_2 - z^5F_2^0$	F
45.890	0		128.10	$a^3P_2 - z^5G_2^0$	H
48.167	6	2W	35100.01	$z^3F_3 - e^3F_4$	F
2857.585	0		34984.33	$a^3P_1 - z^5G_2^0$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
2865.293	0		34890.21	$a^3P_2 - z^5G_3^0$	H
65.378	3	1W	889.18	$z^3F_4^0 - e^5F_4$	H
69.876	6	1W	834.51	$z^3D_3^0 - e^3F_3$	F
73.833	1		785.92	$z^3F_3^0 - e^5F_3$	H
80.134	5	4W	710.45	$z^3D_2^0 - e^3F_2$	F
82.504	0		681.91	$z^3F_2^0 - e^5F_2$	H
83.464	0	2	670.36	$b^3F_3 - z^5F_4^0$	F
2885.165		0	649.9	$b^3F_2 - z^5F_3^0$	H
2918.387	15	?	255.50	$z^3D_1^0 - e^3F_2$	H
22.672	1		205.78	$z^3F_3^0 - e^5F_4$	H
30.239	12	20w	34116.95	$z^3D_2^0 - e^3F_3$	F
42.968	25	15W	33969.40	$z^3D_3^0 - e^3F_4$	F
49.744		0	891.4	$a^5P_1 - z^3D_2^0$	H
51.396	0		872.40	$a^3D_1 - y^5D_0^0$	H
59.589	0w	5	778.63	$a^3P_2 - z^5D_3^0$	H
60.795		0	764.9	$a^3D_1 - y^5D_1^0$	H
61.548	1	0w	756.29	$a^3D_3 - y^5D_3^0$	H
2968.688	0		33675.10	$a^3D_1 - y^5D_2^0$	H
3130.335		0	31936.2	$b^3P_2 - z^5S_2^0$	H
3131.329		0w	926.1	$a^5P_3 - z^3F_4^0$	H
3212.019		3	31124.1	$c^3P_2 - z^5S_2^0$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
3290.761		0	31379.4	$a^5P_3 - z^5G_3^0$	H
3328.380		0	30036.0	$a^5P_3 - z^5G_4^0$	H
52.786		30	29817.4	$a^5P_2 - z^5D_1^0$	M
58.608		10	765.7	$a^5P_3 - z^5D_2^0$	M
70.933		50	656.9	$a^5P_1 - z^5D_0^0$	M
87.698		60	510.1	$a^5P_1 - z^5D_1^0$	M
3388.18		35	505.8	$a^5P_2 - z^5D_2^0$	M
3415.785		75	267.8	$a^5P_3 - z^5D_3^0$	M
23.846		40	198.6	$a^5P_1 - z^5D_2^0$	M
3446.401		50	29007.5	$a^5P_2 - z^5D_3^0$	M
3501.708		200	28549.4	$a^5P_3 - z^5D_4^0$	M
14.219		2	447.7	$a^5P_1 - z^5F_1^0$	M
17.503		10	421.2	$a^5P_2 - z^5F_2^0$	M
17.988		0	417.2	$b^3P_1 - z^3D_1^0$	H
35.940		15	272.9	$b^3P_0 - z^3D_1^0$	H
45.032		25	200.5	$a^5P_3 - z^5F_3^0$	M
55.925		6	114.1	$a^5P_1 - z^5F_2^0$	M
61.123		50	073.0	$a^3G_5 - z^3F_4^0$	H
66.977		5	28026.9	$a^3G_4 - z^3G_3^0$	H
3578.022		30	27940.5	$a^5P_2 - z^5F_3^0$	M
3610.60		3	27688.4	$a^3G_4 - z^3F_4^0$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
3621.177		100	27607.5	$a^5P_3 - z^5F_4^0$	M
43.589		6	437.7	$b^3P_2 - z^3D_3^0$	H
56.18		2	343.2	$c^3P_2 - z^3D_2^0$	H
3656.845		2	27338.2	$a^3G_4 - z^3G_4^0$	H
3714.709		3	26912.4	$c^3P_1 - z^3D_2^0$	H
3754.700		30	26625.8	$c^3P_2 - z^3D_3^0$	H
3983.040		10	25099.4	$a^3D_1 - z^3D_1^0$	H
4040.057		3?	24745.16	$a^3D_3 - z^3D_2^0$	H
4145.153		15	117.8	$a^3D_2 - z^3D_2^0$	H
4160.672		50	24027.8	$a^3D_3 - z^3D_3^0$	H
4272.202		3	23400.6	$a^3D_2 - z^3D_3^0$	H
4285.36		1w	23329.2	$a^3D_1 - z^3F_2^0$	H
4384.25		2?	22802.5	$a^3D_2 - z^3F_2^0$	H
4413.864		30	649.5	$a^5D_4 - z^5P_3^0$	H
82.707		5w	301.7	$a^5D_3 - y^3D_2^0$	H
4489.12		2	269.9	$a^3D_2 - z^3F_3^0$	H
4516.432		10	135.2	$a^3D_3 - z^3F_4^0$	H
33.206		10	053.3	$a^5D_3 - z^5P_2^0$	H
37.971		10	22030.1	$a^5D_3 - z^5P_3^0$	H
59.28		8	21927.2	$a^5D_2 - z^5P_1^0$	H
4569.27		25	21879.2	$a^3H_5 - z^3G_4^0$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
4616.234		5	21656.6	$a^5D_1 - z^5P_1^0$	H
22.011		2w	629.6	$a^5D_2 - z^5P_2^0$	H
60.66		25	450.2	$a^3H_6 - z^3G_5^0$	H
4680.72		0	21358.3	$a^5D_1 - z^5P_2^0$	H
4831.16		10w	20693.2	$a^5D_4 - y^5D_4^0$	H
4865.83		1w	545.8	$a^5D_4 - y^5D_3^0$	H
4980.113		3	20074.3	$a^5D_3 - y^5D_4^0$	H
5012.082		5	19946.2	$a^5D_3 - y^5D_2^0$	H
17.025		3	926.6	$a^5D_3 - y^5D_3^0$	H
5097.445		1	612.2	$a^5D_2 - y^5D_1^0$	H
5120.957		3	522.2	$a^5D_2 - y^5D_2^0$	H
40.114		3	449.4	$a^5D_1 - y^5D_0^0$	H
5192.947	0w	3	19251.50	$a^5D_1 - y^5D_2^0$	H
5477.588	4w		18251.16	$e^5F_5 - v^5D_4^0$	H
5501.548	4w		18171.67	$e^5F_4 - v^5D_3^0$	H
5683.165	1w		17590.96	$e^5F_3 - v^5D_3^0$	H
5773.944	0w		17314.39	$e^5F_5 - w^5D_4^0$	H
5934.252	1 w		16846.66	$e^5F_3 - w^5D_3^0$	H
6083.964	3w		432.11	$e^5F_4 - x^5F_3^0$	H
86.496	1w		425.28	$e^5F_5 - x^5F_4^0?$	H
6088.017 .	1w		16421.17	$^3\gamma_3^0 - e^3F_4$	H

Table XIV. (Continued)

WAVE LENGTH IN AIR	S.T. INT.	SPK. INT.	WAVE NUMBER IN VACUUM	COMBINATION	AUTHOR
6201.034	lw		16121.89	$e^5F_3 - x^5F_2^o$	H
65.561	lw		15955.86	$e^5F_5 - x^5F_5^o$	H
6292.156	ow		888.42	$^310_2^o - e^3F_3$	H
6305.768	lw		854.12	$e^5F_4 - x^5F_4^o$	H
06.655	ow		851.89	$e^5F_3 - x^5F_3^o$	H
6370.218	lw		15693.72	$e^5F_2 - x^5F_2^o$	H

VIII. CONCLUSION

Although a considerable majority of the strong lines in the spectrum have now been classified, many still remain unidentified. Most of these should be the result of combinations between the many odd triplets of the $3d^7 4p$ configuration and the even triplets of the $3d^7 4s$ configuration. At least traces of some of these triplets have been found, but it has not been possible to completely establish them, even though it is felt that their approximate positions are well enough known so that it should not have been difficult to establish them. Apparently neither the condensed spark nor the Schuler tube provide the proper excitation to permit completion of this phase of the analysis. In order to further develop these lower levels of the spectrum, it is recommended that the method of the cathode spot be tried. The work of Dobbie^{16a} on the Fe II spectrum shows that there is some hope for success with this type of excitation.

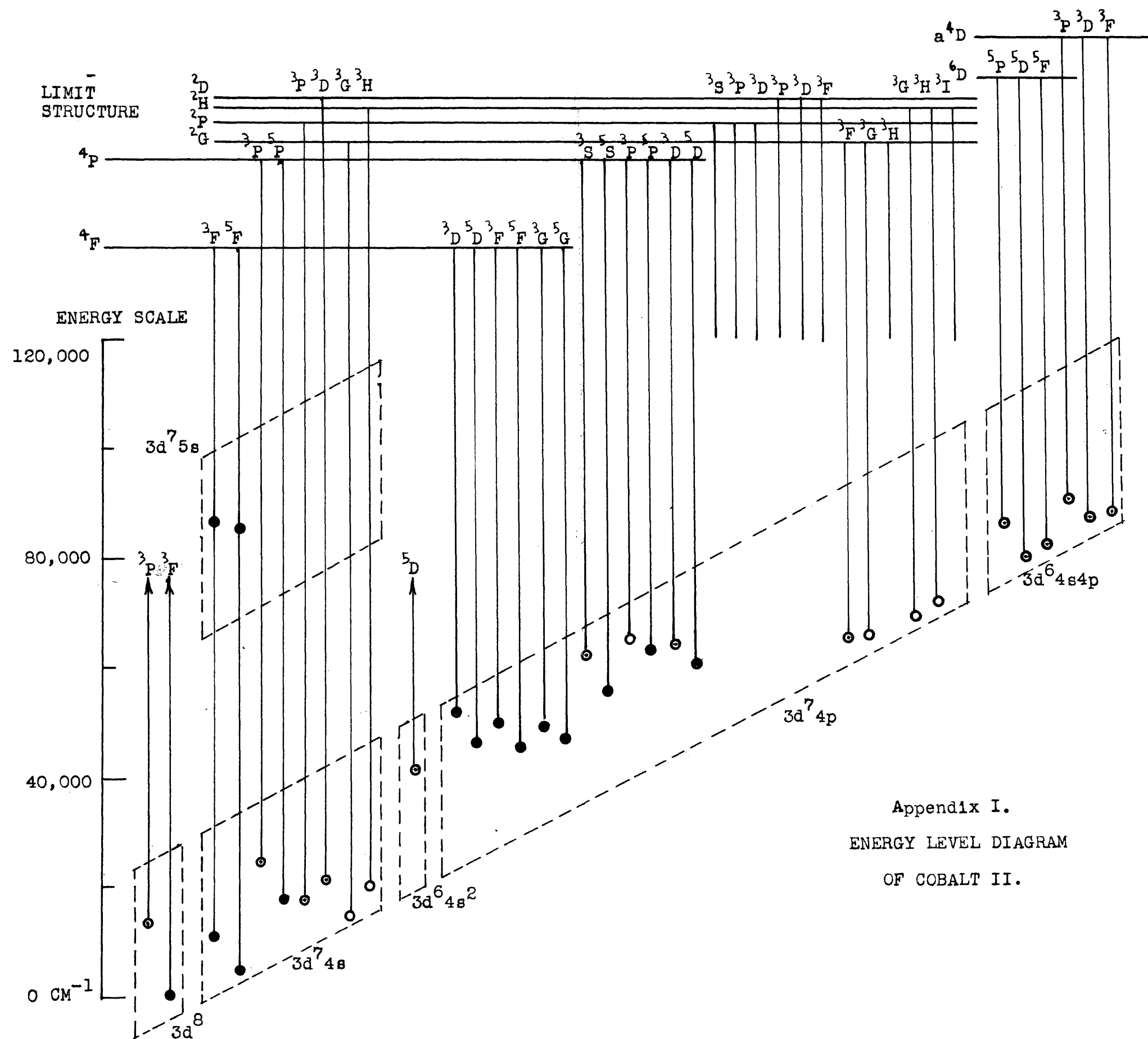
It now appears that the Schuler tube, employed in the manner described in Appendix II, is not a useful source for extending the study of the Co II spectrum in the

16a - J. C. Dobbie, Annals of the Solar Physics Observatory, Volume V, Part I, Cambridge (1938)

region of higher energy levels. For further development of these high levels, there is need for some new method of excitation. It should be stated that the experimental work was not completely fruitless. Most of the lines previously observed in the condensed spark were also observed in the hollow cathode discharge, and because of the great sharpness of the lines produced by the Schuler tube, and because of the high dispersion of the 21 ft. grating instrument, the measurement of these lines has been accurate. Consequently the energy levels of Co II have been established with a higher degree of accuracy than heretofore. Also the terms e^3F and e^5F in the $3d^75s$ configuration have been given much support by the appearance of 14 new lines which did appear in the discharge.

KEY TO SYMBOLS

- Designates terms found by Meggers and Findlay
- ⊙ Designates new terms established by Hager
- Designates incomplete or tentatively established new terms



Appendix I.
ENERGY LEVEL DIAGRAM
OF COBALT II.

Appendix II.

SCHULER TUBE DEVELOPMENT

The condensed spark is unsatisfactory for exciting the high levels of Co II. The few lines which do appear as a result of combinations with high levels cannot be measured accurately, because they are broad, asymmetric, and very weak. It is necessary to employ an excitation agent other than the spark to establish the high series members required for more accurately determining the ionization potential. A possible source is the hollow cathode discharge tube, called the Schuler tube, which has been useful for exciting high levels in the spectra of several similar metallic ions.

The work of Paschen, Schuler, Shenstone and others has shown that when a hollow cathode discharge is run in a noble gas, and when a sample of a metal is placed inside the cathode, the spectrum of the metal will be produced, and marked enhancements of certain lines will appear. Sawyer¹⁷ suggests that free atoms of the metal experience collisions of the second kind with noble gas atoms which have either been ionized or excited to metastable states. As a consequence of the collision, the ionization energy or the excitation energy may be transferred from the noble

17 - R.A. Sawyer, Phys. Rev., 36, 44 (1930)

gas atoms to the metallic atoms. One may observe enhancements of those lines in the spectrum of the metallic atom which originate in levels whose energy values, as measured above the ground state of the neutral atom, correspond to the excitation energy obtained from the atoms of the gas.

In the case of many metals, the levels where maximum enhancements occur fall somewhere in the spectrum of the singly ionized atom, depending on the noble gas used in the tube. By prudently choosing gases, it is possible to achieve selective excitation in several energy level regions. Table XV lists the levels at which three of the noble gases may produce enhancements in the spectrum of singly ionized cobalt. The energy required to ionize the cobalt atom must, of course, be accounted for when computing the values of these levels. The table shows that both neon and helium should be useful for exciting those upper series members which are predicted to fall in a region between 80,000 and 100,000 wave numbers above the ground state of the singly ionized cobalt atom.

The discharge must not only maintain the collision process which produces the excitation, but it must also provide the means for freeing metallic ions from the solid

Table XV.

ENHANCEMENT BY NOBLE GASES

GAS	IONIZATION POTENTIAL CM-1	METASTABLE STATES CM-1	IONIZATION POTENTIAL CM-1	LEVEL OF ENHANCEMENT CM-1
He	198308		63400	134,908
		159830	63400	96,430
		166251	63400	102,851
Ne	173930		63400	110,530
		134819	63400	71,419
A	127104		63400	63,704
		94546	63400	31,146

body of the sample inserted in the cathode. It does this by cathodic sputtering or boiling, and cobalt cannot be expected to respond readily to either process, because it will boil only at very high temperatures as shown in Table XVI, and it is not a good sputterer, as shown in Table XVII.

Table XVI.

MELTING AND BOILING POINTS OF METALS¹⁸

METAL	MELTING POINT	BOILING POINT
Co	1495 °C.	3000 °C.
Fe	1535	3000
Cu	1083	2300
Pt	1774	4300
Mn	1260	1900

18 - Handbook of Chemistry and Physics, p. 1898 (1950-51)

Table XVII.
SPUTTERING RATES OF METALS¹⁹

METAL	SPUTTERING RATE (mg/amp-hr)
Bi	1470
Te	1200
As	1100
Tl	1080
Sb	890
Ag	740
Au	460
Pb	400
Zn	340
Cu	300
C	262
Sn	196
Fe	68
Ni	65
W	57
Co	56
Mo	56
Mn	38
Cd	32
Al	29
Cr	27
Ta	16
Mg	9

- 19 - Gunterschulze, Zeits fur Physik, 36, 563 (1926)
These measurements were made in hydrogen with a
770 volt cathode fall, and a current density of
7 milliamperes per square centimeter.

Shenstone succeeded in running the Schuler tube with Cu²⁰ and Pt²¹. Although platinum is seen to have a very high boiling point, it is a moderately good sputterer²². Curtis²³ was successful in exciting Mn by this method, and in this case it is likely that the successful excitation may have been due to the relatively low boiling point of the metal, because the sputtering rate of this metal is quite low. Green²⁴ had difficulty in exciting iron, and cobalt is unfortunately very similar to iron with respect to both boiling point and sputtering characteristics.

20 - A. G. Shenstone, Phil. Trans. Roy. Soc. 237
453 (1938)

21 - A. G. Shenstone, Phil. Trans. Roy. Soc. 235
195 (1936)

22- Platinum is not included in Table XVII, but is considered to be as good a sputterer as copper by other observers. See, for example, J. Strong, "Procedures in Experimental Physics", p. 163, Prentice-Hall, (1938)

23 - C. W. Curtis, Phys. Rev. 53, 474 (1938)

24 - L. C. Green, Thesis, Princeton University (1937)

A. THE POWER SUPPLY

Because the Schuler tube has negative resistance characteristics, i.e., the potential drop across the tube decreases when the current increases, it was necessary to provide some external means for stabilizing the current. The simple series ballast resistor is not practical in this case because a potential of 1000 volts is required to pass a current of more than 2 amperes through the Schuler tube, and the potential drop across such a ballast resistor would make it necessary for the power supply to furnish a potential considerably higher than that required by the Schuler tube itself. Increasing the voltage of the power supply is very costly when such large currents are required.

The power supply used was a three phase constant current source designed by Green²⁵ for running the Schuler tube in the iron experiment. This system was partially re-designed to permit use of easily obtained components. The circuit is shown in Fig. 4, and the specifications for the parts are listed in Table XVIII. A mathematical analysis of the circuit is given in the above reference.

Fig. 5 shows how nearly constant the current remains when the load resistance is varied. The safety cutoff circuit is mandatory, because an open load circuit pro-

25- L. C. Green and J. B. Horner Kuper, Rev. Sci. Instruments, 11, 250 (1940)

Fig. 4

POWER SUPPLY CIRCUIT

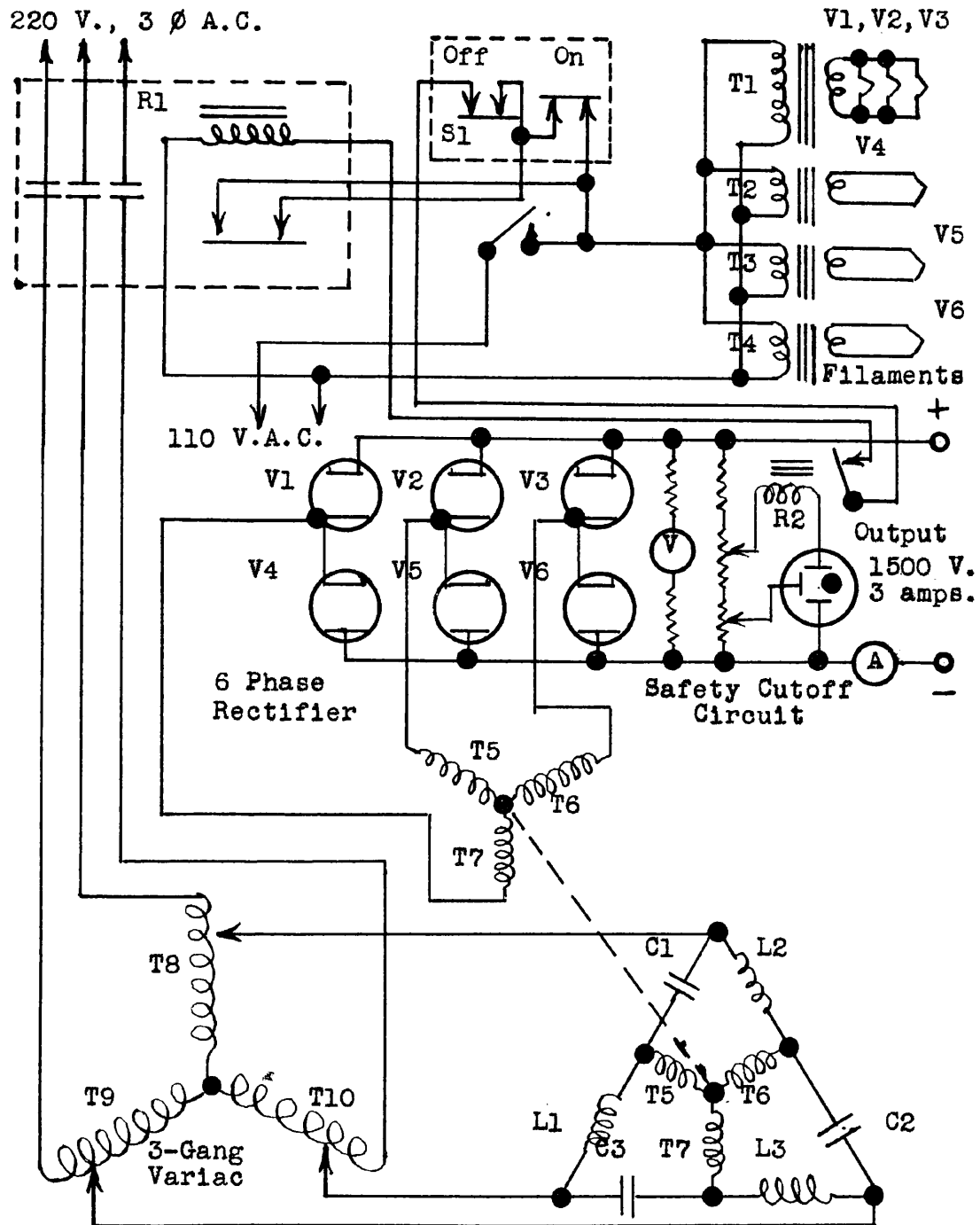


Fig. 5
CONSTANT CURRENT POWER SUPPLY

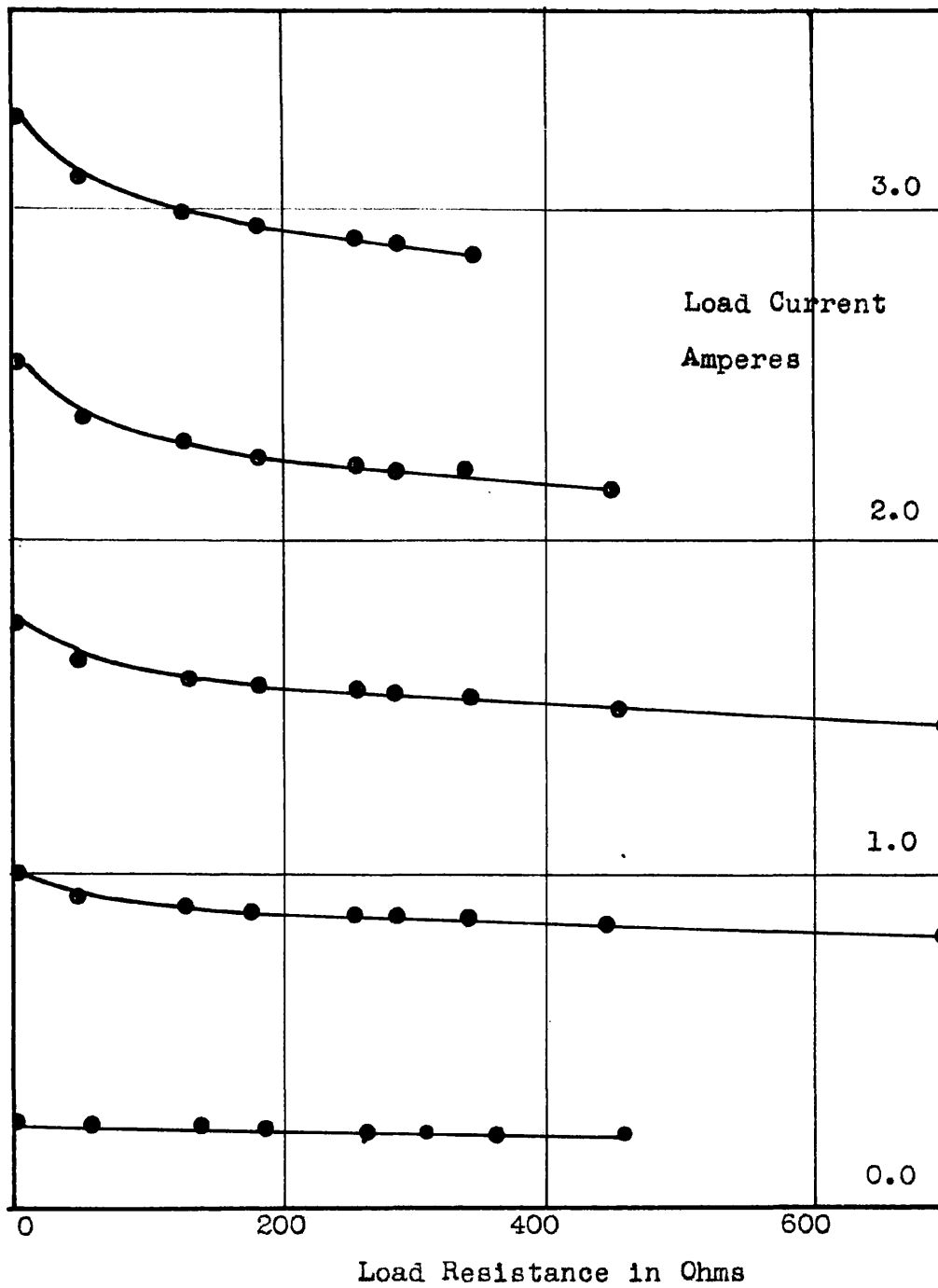


Table XVIII.

SPECIFICATIONS FOR CONSTANT CURRENT POWER SUPPLY

DESIGNATION IN FIG. 4	DESCRIPTION
V1 - V6	Type 872-A mercury rectifier tubes
V7	Type OA4G cold cathode thyatron
T2, T3, T4	Filament transformers, 5 v., 10 amps.
T1	Filament transformer, 5 v., 30 amps.
T5, T6, T7	Standard $1\frac{1}{2}$ K.V.A. stepdown transformers, connected so that ratio of turns is 4:1, and used backwards to step voltage up.
T8, T9, T10	Standard Variacs, 115 v., 10 amps., ganged on single control shaft. (115 volt Variacs run safely on 230 volt 3 phase line, if they are connected in a wye network.)
R2	Sensitive plate circuit relay.
R1	Standard motor start relay with lock-in contacts.
S1	Spring-load push switches
S2	Single pole single throw switch
C1, C2, C3	Banks of condensers totalling 125 mfd. each. Should be oil capacitors rated at 330 V. A.C.
L1, L2, L3	Inductances made of #8 D.C.C. wire wound on a hollow-cored spool. The core is 10 cm. in diameter and 15 cm. long. The outside diamter of the wound coil is 19 cm. The value of the inductance was adjusted to give resonance with the capacitors by inserting iron wire into the open core while passing 60 cycle current through a series combination of the inductance and the capacitance and adjusting the voltage applied to the combination to keep the current at about 15 amperes.
Resistors chosen to match relays and meters.	

duces a large voltage drop across the large banks of condensers in the constant current network, and these condensers may be completely destroyed.

B. THE VACUUM SYSTEM

Fig. 6 shows the vacuum system employed in conjunction with the Schuler tube. It had two functions: with stopcocks A, E, and F open, and with stopcocks B, C, and D closed, the system served to evacuate the Schuler tube and the associated traps. With A, C, and D closed and B, E, and F open, the diffusion pump circulated the noble gas through the system. The gas was inserted by opening and closing D, putting gas into the small dosing chamber, and then opening and closing C. Liquid air was used on the traps.

A "getter" tube using a hollow cathode discharge over mischmetal²⁶ was provided for cleaning small impurities

26 - Mischmetal may be obtained from New Process Metals Co., 46 Center St., Newark, N.J. Airco "reagent grade" noble gases were found to be pure enough so that the tube ran well without mischmetal cleaning, provided the cathode was well out-gassed, using an induction furnace, or by heating the cathode in the Schuler tube discharge and then pumping out the gaseous impurities from the system before dosing with clean noble gas.

Fig. 6
THE VACUUM SYSTEM

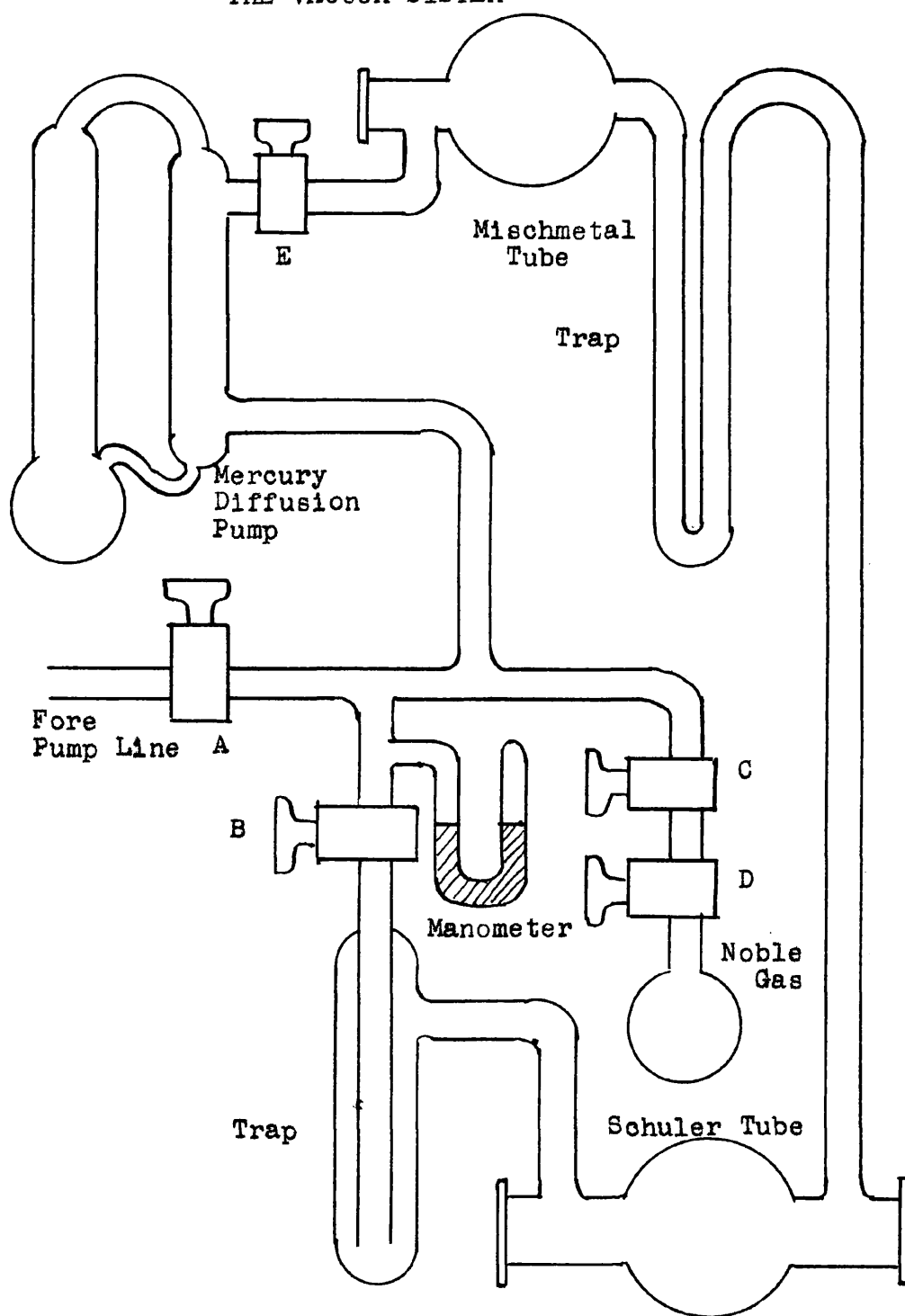
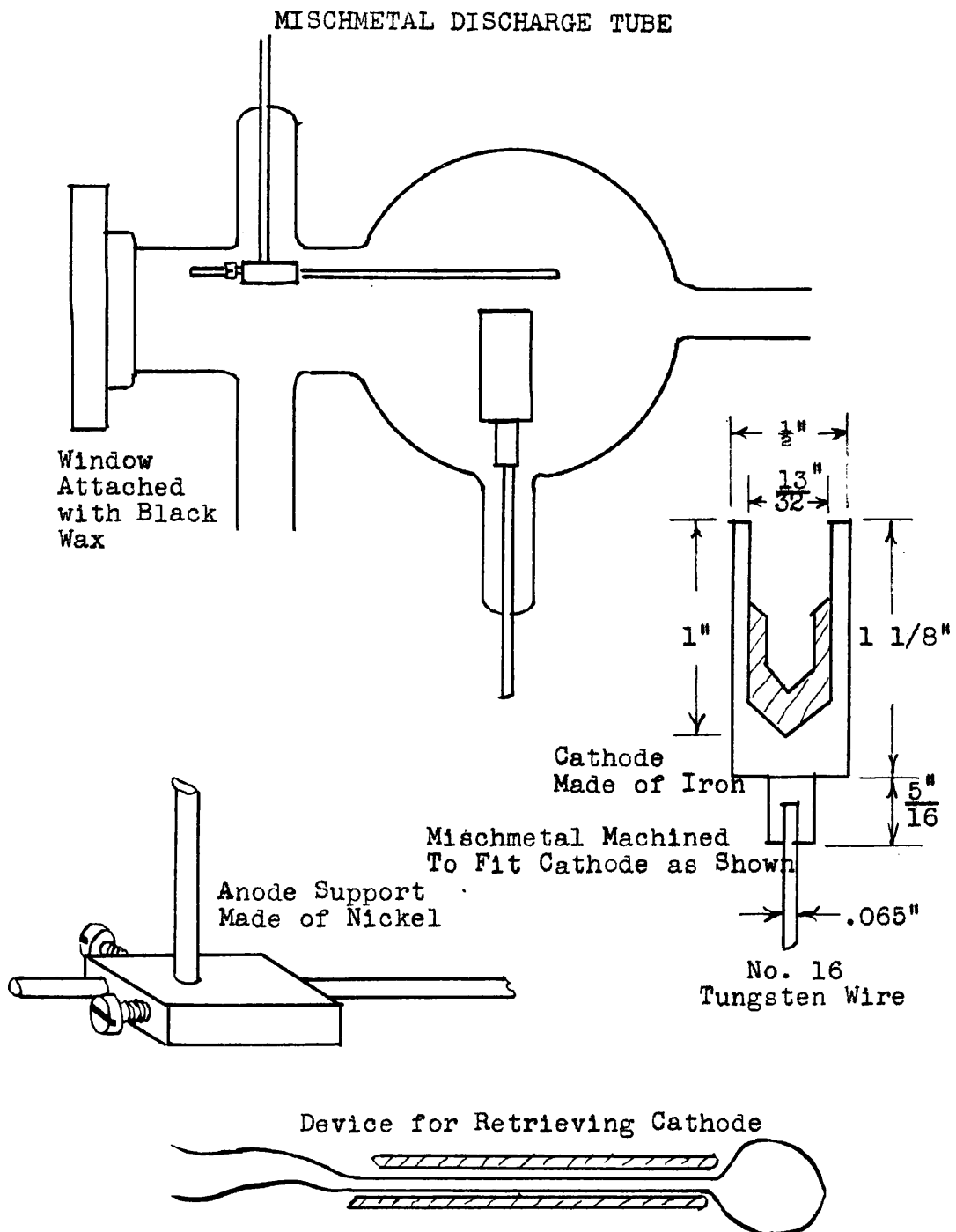


Fig. 7



from the gas in the system. Fig. 7 shows the construction details of a mischmetal discharge tube with a cathode that is removable for re-charging. After the window and the anode were removed, the cathode was lifted out quite easily with a section of glass tubing and a piece of wire arranged as shown in the figure. When operated at a pressure that produced a 5 mm. differential on the mercury manometer, the mischmetal discharge required a current of 0.25 ampere and 3 to 4 hundred volts potential across the tube.

C. THE SCHULER TUBE

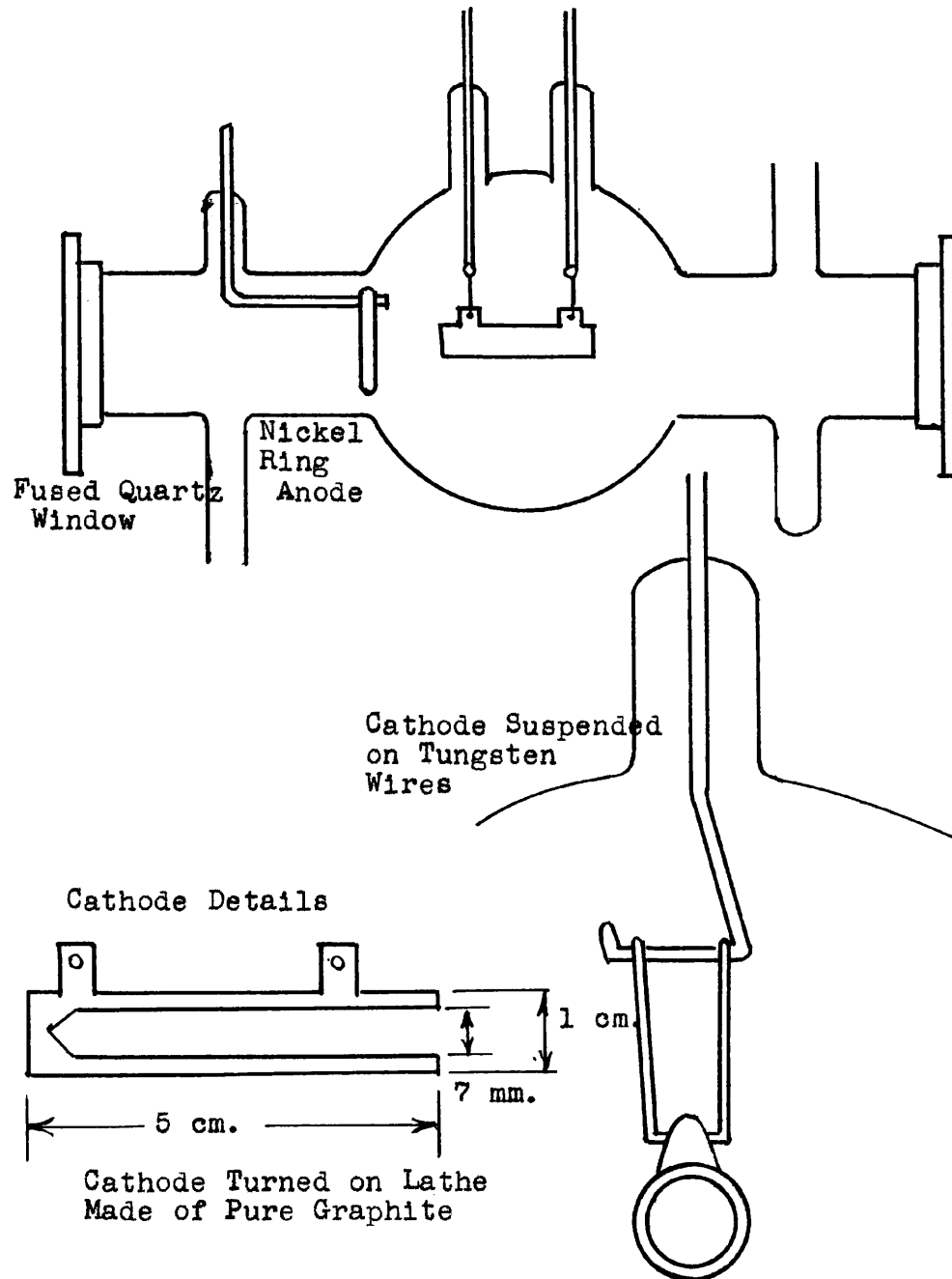
The Schuler tube used was a 3-liter Pyrex flask with a second neck fastened on the side of the bulb opposite the original neck, as shown in Fig. 8. The hollow cathode was machined from graphite²⁷ on a lathe. The inside diameter of the most satisfactory cathode was 7 millimeters.

The cathode was loaded with about five 1 mm. cubes of pure cobalt and was suspended inside the tube with tungsten wires. The windows were fastened onto both ends of the tube with black wax. An electric fan circulated air around the tube during running to prevent softening the wax and overheating of the main bulb.

27 - ACHESON GRAPHITE rod, 1 inch in diameter, was obtained from the National Carbon Div. of Union Carbide and Carbon Corp., 30 E. 42nd St., New York.

Fig. 8

DETAILS OF THE SCHULER TUBE



Before running the discharge, the system was evacuated to about 10^{-5} mm. of mercury, and it was then dosed with noble gas until the manometer showed a differential of a few millimeters. With the system used, this required about two doses, with a 3 c.c. dosing chamber, when the supply flask was new, i.e., when the noble gas pressure in the flask was approximately equal to atmospheric pressure.

Then the constant current source was turned on at low power, and it was gradually turned up. If the discharge did not take place inside the cathode, the pressure was increased until it did. If, on the other hand, there was excessive arcing, the pressure was decreased.

When the current reached about 2 amperes, the cathode was white hot. Usually, the cobalt spectrum could be observed when the current was above 1.5 amperes. Often, when there was a fresh charge of cobalt inside the cathode, excessive arcing occurred which could not be eliminated by lowering the pressure. But after running for a few minutes at 1.5 amperes or more, the discharge settled down and ran smoothly. After about 30 minutes of running, the suspending wires began to heat up slowly. After another 30 minutes, the wires became so hot that the tube had to be turned off.

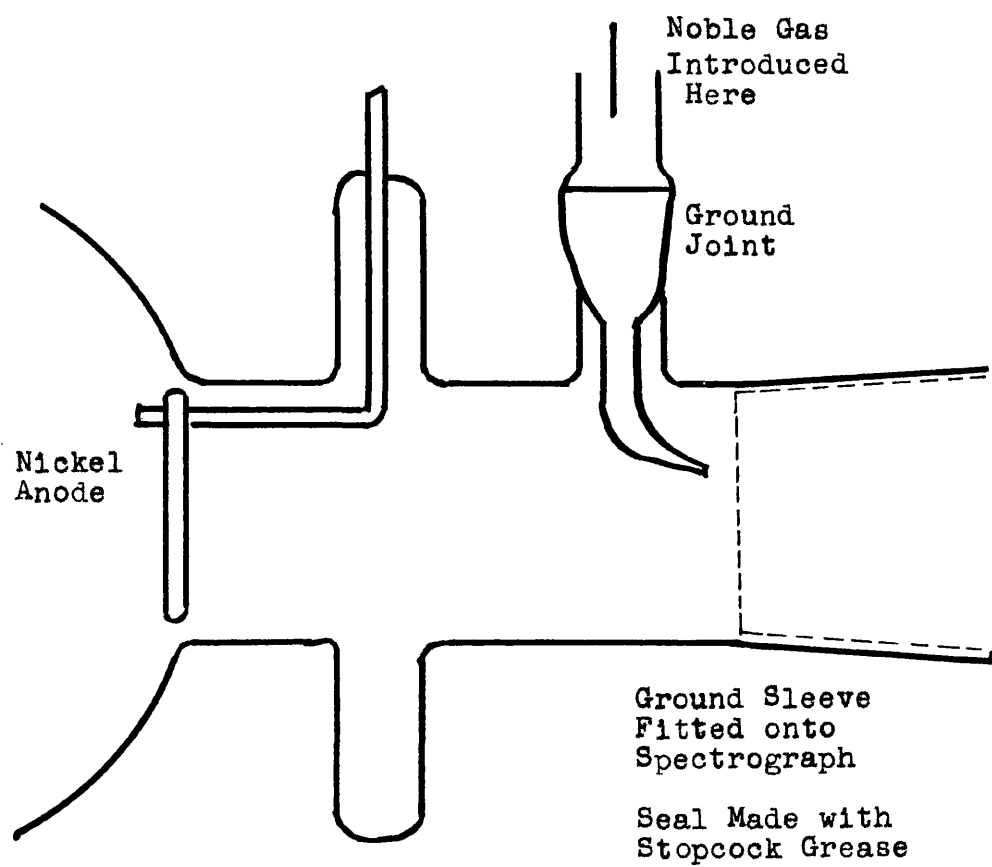
No change in the geometry of the cathode or the method of suspending the cathode permanently removed this trouble, but a small fresh lump of cobalt in the cathode always returned the tube to proper running condition for another 30 minute period.

When the Schuler tube was used with the vacuum spectrograph, only a few changes had to be made. They are shown in Fig. 9. Instead of circulating the gas as before, the gas was injected into the tube through a glass nozzle which directed the stream of gas toward the slit of the spectrograph. The gas passed through the slit into the vacuum chamber. This served to charge the tube with the noble gas and also to keep the slit clear of matter that was thrown from the cathode from time to time during the running of the discharge. The noble gas used was helium, and it was purified with a chabazite²⁸ trap. The rate of flow into the Schuler tube was kept uniform by passing the gas through a small artificial leak.

28 - Purification by chabazite is described by
G. P. Baxter and H. W. Starkweather, Pro.
Nat. Acad. Sci., 11, 231 (1925)

Fig. 9

SCHULER TUBE MODIFIED FOR VACUUM SPECTROGRAPH



D. THE EXPERIMENTAL RESULTS

The results of the experiment were not gratifying in that no new terms have been established as a result of the plates obtained with the Schuler tube.

It is true that some atoms of cobalt must have been released from the solid sample because the cobalt spectrum was produced. It is also evident that some of the desired collision excitation was occurring because so many lines appeared that resulted from combinations with the high terms e^3F and e^5F . These lines were very sharp and could not have resulted from the type of excitation produced by either the arc or the spark. But it seems that the collision excitation did not take place on a large enough scale to produce the lines needed to fix the position of new energy levels. The extremely high temperature required to produce free cobalt atoms may have been a factor in reducing the excitation due to collisions of the second kind. Also, at this high temperature, molecular band spectra appeared, causing some sections of the plates to be so over-populated with lines that they were virtually unmeasurable.

With neon in the Schuler tube, the spectrum was photographed from 2000 to 6400 A. The plates were measured from 2000 A to 3000 A, and from 5000 to 6400 A. The region from 3000 to 5000 A did not seem worth measuring, especially when it became clear that the enhancement was poor. With helium in the tube, a plate was exposed on the vacuum spectrograph. This plate was measured from 1400 to 1900 A, and again the enhancement was poor.

Appendix III.

MASTER LINE LIST

The following pages are photostatic reproductions of a master line list which is a compilation of all of the available measurements of the Co II spectrum.

The wave number values calculated from wave length measurements made of the Schuler tube plates are listed with one more significant figure than the values reduced from spark plates.

The intensity estimates for lines taken from the Schuler tube plates are based on a 0-200 scale. In the wave length region 1010 to 2225 Å the intensity of the spark lines is estimated on the basis of a 1-16 scale, except where the estimates are followed by the letters "F" and "M", in which case they are based on a 0-200 scale. In the region from 2225 to 5140 Å the spark intensities represent the average of several estimates, and the 0-200 scale was used.

λ vac	λ_{ST}	λ	λ vac	λ_{ST}	λ	λ vac	λ_{ST}	λ	λ vac	λ_{ST}	λ
1664.775	3	60068.2	1725.679	3	57948.2	1826.394	4	54752.7	1962.486	1	50955.8
66.276	119	014.1	$\alpha^3D_3 - \omega^3D_3^0$	28	910.5	27.036	1	731.7	63.537	1	928.5
66.380	5	60010.3	26.802	1	976.9	27.919	1	707.0	64.810	1	914.2
68.613	2	59930.0	27.899	0	807.2	28.274	1	696.4	68.692	3	795.1
69.907	4	868.4	30.815	2	776.3	28.409	2	672.4	68.680	5	769.7
70.334	1	2	868.3	1	711.3	29.397	1	662.8	70.731	2	742.6
70.799	0	2	851.6	0	692.3	32.471	0	571.1	71.158	1	731.6
70.943	0	1	846.4	4	644.0	33.271	1	547.3	72.530	2	696.3
71.553	12	824.6	34.125	0	651.8	33.475	3	541.2	73.855	1	662.3
73.145	2	3	767.7	0	629.9	33.445	1	535.6	75.008	4	632.7
73.787	1	1	744.8	0	623.9	34.251	1	518.2	75.447	2	614.3
73.918	3	740.1	$\alpha^3D_3 - \alpha^3F_3^0$	2	596.2	35.940	0	468.0	77.304	1	573.9
74.921	1	5	704.3	0	514.5	36.447	6	447.0	77.549	4	567.6
75.183	2	695.0	$\alpha^3D_3 - \alpha^3F_3^0$	10	509.4	37.556	0	420.1	78.334	1	547.6
75.758	1	2	674.5	0	488.3	37.993	6	407.2	80.251	1	498.6
77.000	0	1	630.3	3	480.9	38.844	2	382.0	80.415	3	494.5
77.093	0	1	627.0	1	452.3	39.263	4	369.6	80.854	1	483.3
77.473	0	1	613.5	0	424.4	39.356	12	366.9	81.601	2	464.2
77.685	1	606.0	41.874	2	408.8	40.809	6	341.7	81.736	2	460.8
79.232	2	551.0	$\alpha^3D_3 - \alpha^3F_3^0$	0	359.8	42.350	4	278.5	82.259	3	447.5
79.368	3	546.2	43.571	1	353.6	42.811	9	264.9	82.530	1	440.6
79.490	2	541.9	43.791	1	344.3	44.009	0	229.7	83.753	4	404.4
79.558	2	539.5	46.297	2	264.0	45.106	1	197.4	84.843	3	381.8
80.125	3	519.4	47.119	10	237.1	45.916	5	173.6	84.930	4	324.4
80.285	1	513.7	47.882	5	212.1	48.459	8	099.1	87.651	1	310.6
80.897	0	492.0	48.343	3	197.0	48.814	0	057.5	89.429	1	265.7
81.165	0	482.6	48.613	1	188.2	51.641	4	54806.1	89.803	1	256.2
82.102	10	3	449.4	0	183.0	52.553	3	53780.3	90.356	1	242.3
82.821	5	6	424.0	0	144.1	52.693	3	976.0	91.567	3	211.7
83.164	2	411.9	54.217	7	7005.5	53.715	2	945.7	93.880	3	153.5
83.722	3	6	392.2	0	5690.8	55.037	2	907.3	94.502	1	135.3
84.400	0	1	354.2	2	870.4	56.144	1	875.1	95.733	2	106.7
85.251	1	338.3	58.442	6	862.1	59.625	2	774.3	97.135	3	871.7
85.848	3	317.3	59.085	1	847.7	60.209	5	757.4	1998.549	1	50836.3
86.600	1	290.9	59.834	1	823.5	61.420	1	722.4			
86.896	1	5	280.5	1	764.1	61.855	8	709.9			
87.057	1	5	274.8	0	757.5	62.353	6	695.5			
87.596	0	255.9	61.056	2	751.9	63.842	0	652.0			
87.768	1	249.1	63.475	5	699.8	65.868	0	594.4			
87.848	2	247.0	63.889	0	692.9	67.008	1	561.6			
88.296	2	3	231.3	4	680.4	68.322	1	524.0			
88.681	1	2	217.8	0	419.0	69.211	8	498.5			
89.027	3	205.7	73.040	1	400.3	71.770	0	425.4			
91.310	2	4	125.8	0	373.9	74.150	2	367.5			
92.650	2	079.0	75.867	3	310.5	75.726	2	312.7			
93.330	1	3	055.2	3	253.8	77.332	1	267.1			
93.393	3	053.0	76.318	4	232.9	85.530	1	035.5			
93.816	1	5903.3	81.068	1	146.1	86.664	0	53003.6			
94.072	0	58029.4	81.553	0	130.8	90.016	2	51707.9			
95.894	1	765.7	83.909	0	056.7	91.028	0	861.3			
96.962	3	92.8.8	84.801	0	56028.7	91.919	0	856.4			
97.630	0	905.7	87.349	3	55948.8	95.477	0	757.2			
98.126	1	588.4	87.952	2	927.9	189647.4	1	729.4			
1699.343	0	846.3	88.458	0	914.1	190039.9	0	620.5			
1700.487	1	806.7	93.499	1	756.7	03600	2	532.1			
00.621	0	802.0	93.920	1	743.8	08518	0	396.7	$\alpha^3P_2 - \alpha^3P_1^0$		
02.403	1	733.6	96.608	1	691.4	08738	0	390.6			
02.916	0	722.8	96.660	12	658.8	12563	0	285.9			
03.821	1	691.6	98.786	1	593.0	13756	2	253.3	$\alpha^3P_1 - \alpha^3P_1^0$		
05.643	0	629.0	99.271	0	578.1	17627	1	147.8	$\alpha^3P_2 - \alpha^3P_2^0$		
06.062	3	614.5	99.554	0	569.3	18747	10	111.9			
06.794	1	589.4	102.128	1	490.0	20706	3	064.2	$\alpha^3P_0 - \alpha^3P_1^0$		
06.878	1	586.5	1804.947	8	403.3	21457	0	043.8			
07.563	2	563.0	07.267	1	382.2	22561	3	014.2			
07.775	1	558.7	07.956	0	311.1	22927	5	52004.1	$\alpha^3P_1 - \alpha^3P_2^0$		
09.553	2	495.5	08.276	7	301.3	24526	0	51960.8			
11.101	3	444.1	08.912	6	281.8	27040	4	873.1			
12.128	1	406.	12.167	2	182.6	28498	0	853.8			
12.373	1	398.5	13.138	1	153.0	28956	1	841.5			
12.614	1	390.3	15.330	0	086.4	29388	0	829.9			
13.379	10	364.2	16.608	1	047.6	33.624	1	716.4			
13.426	2	362.6	16.881	1	039.4	35.963	1	663.9	$\alpha^3P_2 - \alpha^3P_2^0$		
15.053	2	307.2	17.459	1	55021.9	36571	3	837.7			
15.873	1	279.4	18.315	1	54996.0	36770	1	632.4			
16.163	1	269.5	18.542	2	989.1	37849	4	624.4	$\alpha^3P_1 - \alpha^3P_1^0$		
17.104	3	237.5	18.649	1	985.9	37394	1	615.7			
18.280	0	197.7	20.310	4	935.7	40.64	0	512.3	$\alpha^3P_2 - \alpha^3P_2^0$		
18.464	0	191.5	21.595	2	896.9	44.186	3	435.4	$\alpha^3P_1 - \alpha^3P_1^0$		
20.168	2	133.9	21.793	3	890.9	58.098	0	279.5	$\alpha^3F_2 - \alpha^3D_2^0$		
23.013	4	037.9	23.891	0	827.8	58.424	1	270.9			
23.324	0	827.4	24.849	1	823.1	51.887	2	232.4			
23.726	0	013.9	25.328	8	764.7	53.071	3	201.4			
23.981	2	58005.3	25.657	1	774.8	55.148	2	149.6			
1725.223	2	57963.5	1826.098	3	54761.6	19574.24	0	087.6	$\alpha^3F_2 - \alpha^3D_2^0$		

60068

106

50036

[illegible]

[illegible]

λ	Wt	λ
6267.616	0w	15950.63
63454	0w	945.95
63945	0w	944.70
70.877	2w	942.33
71.076	2w	941.83
72.746	1w	937.58
74.914	0w	931.08
75.835	0w	929.74
76.350	1w	928.43
79.033	2w	921.63
79.224	0w	921.14
82.976	1w	911.64
83.835	1w	909.46
84.911	0W	906.73
86.321	0w	903.17
88.215	0w	898.38
88.791	2w	897.05
89334	2w	894.03
92.154	0w	888.42 $^310^0 - e^2F_3$
92.742	5w	884.42
95.203	0W	880.73
95.609	0w	879.70
97.139	0w	875.85
6299.456	0w	870.01
6305.268	1w	854.12 $^3F_4 - n^5F_4^0$
86.655	0W	851.89 $^3F_3 - n^5F_3^0$
11.119	1W	840.68
12.440	0w	837.26
21.744	2w	814.06
24.125	0w	808.10
24.490	1w	807.19
26.961	0W	796.02
33.226	1W	785.39
33.419	1w	784.75
35563	0w	779.56
38.831	0w	771.43
46.875	0w	753.42
46.455	0w	752.48
46.580	0w	752.17
47.182	6w	750.88
48.529	1w	747.33
50.872	3S	743.51
51.849	0w	739.11
56.972	0w	726.42
57.707	4w	724.60
62.687	0w	712.29
64.239	10-100	706.47 4I
70.218	1w	693.72 $^3F_3 - n^5F_2^0$
71.241	0w	691.20
71.796	5w	689.84
73.789	30	684.93
84.156	0w	659.46
86.901	6w	652.73
88.189	2S	649.56
89.406	5W	646.59
90.751	3w	643.30
95.779	2W	630.96
6397.148	4	627.58
6401.287	0w	15617.55

Appendix IV.

VITA

Nathaniel Ellmaker Hager, Jr. was born on June 3, 1922 in Lancaster, Pennsylvania, the son of Nathaniel Ellmaker and Ruth Mayer Hager.

He graduated from Franklin and Marshall Academy in June 1940 with honors in science and literature, and then attended Franklin and Marshall College. He graduated in November of 1943 with the degree of Bachelor of Science with Departmental Honors in Physics, and while at Franklin and Marshall he was elected to Phi Beta Kappa and Sigma Pi Sigma honorary fraternities.

He was commissioned as an Ensign in the United States Naval Reserve in January 1944, and served with the United States Submarine Forces until July 1946, when he retired from active duty with the rank of Lieutenant Junior Grade.

He entered Lehigh University in March 1947, and received the degree of Master of Science in June 1948. While at Lehigh he became an associate member of Sigma Xi honorary fraternity. He has served at Lehigh as a graduate assistant and as an instructor in the physics department.

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